# A Quasi-Newton Method with No Derivatives

### By John Greenstadt

Abstract. The Davidon formula and others of the "quasi-Newton" class, which are used in the unconstrained minimization of a function f, provide a (generally) convergent sequence of approximations to the Hessian of f. These formulas, however, require the independent calculation of the gradient of f. In this paper, a set of new formulas is derived—using a previously described variational approach—which successively approximates the gradient as well as the Hessian, and uses only function values. These formulas are incorporated into an algorithm which, although still crude, works quite well for various standard test functions. Extensive numerical results are presented.

1. Introduction. The so-called variable-metric method for minimizing functions, which was discovered by Davidon [1] and developed by Fletcher and Powell [2], has been so successful that it has attracted a great deal of interest. Various theoretical studies, as well as new, related algorithms, have appeared in the literature ([3]-[6], among many others).

So far, all but one\* of these variants of the DFP (Davidon-Fletcher-Powell) method have required the explicit evaluation, at each step, of the gradient of the function f to be minimized. From these computed gradients, the inverse of the Hessian matrix is gradually constructed, and the Newton formula (which is used to compute the next step direction) becomes gradually more accurate.

In a previous publication [7], it was shown how DFP-like formulas could be derived by solving a certain variational problem. In this paper, the same method will be applied to finding quasi-Newton\*\* formulas which do not involve the explicit calculation of gradients. Clearly, since the gradient is needed in the Newton formula, the new algorithm will have to estimate it—as well as the Hessian— in the same way as the inverse Hessian is estimated in the DFP method.\*\*\*

The basic notation to be used is as follows: f(x) is the function of the variables  $(x_1, x_2, \dots, x_N)$  in  $R_N$  which is to be minimized;  $\tilde{g}$  and  $\tilde{G}$  are the gradient and Hessian of f, respectively. In the course of the work, certain estimates of these quantities will be discussed; these will be denoted by g and G (without bars). Further,  $H \equiv G^{-1}$ . At certain stages, vectors specifying directions for line searches are introduced; the letter d is used to denote these. When a direction vector d has been normalized (in a sense to be outlined later), the normalized direction is denoted by the letter s. Using a starting point  $x_0$  and a unit direction s, a straight line in  $R_N$  may be expressed parametrically as follows:

Copyright © 1972, American Mathematical Society

Received July 27, 1970, revised July 26, 1971.

AMS 1970 subject classifications. Primary 90C30.

Key words and phrases. Quasi-Newton methods, gradient-free nonlinear optimization.

<sup>\*</sup> This is the method of Stewart [4] which, however, computes the gradient by finite differences.

**<sup>\*\*</sup>** The term "variable-metric" is reserved by convention for those methods in which the Hessian remains positive-definite (and hence can be regarded as a "metric" tensor).

<sup>\*\*\*</sup> A method due to Fiacco and McCormick [16] also estimates the gradient and Hessian using only function values. A comparison is made in Appendix B.

JOHN GREENSTADT

$$(1.1) x(\alpha) = x_0 + \alpha s,$$

where  $\alpha$  is a parameter which measures the distance from  $x_0$  to  $x(\alpha)$ . If a line search along such a line has terminated at a certain value  $\alpha_1$ , the displacement vector  $x(\alpha_1) - x_0$  will be denoted by  $\sigma$ , so that  $\sigma = \alpha_1 s$ .

At appropriate places, subscripts may be appended to any of these symbols, to label the various steps with which they are associated. At other places, the context permitting, the subscripts will be dropped.

2. The Role of the Constraint in the DFP Case. In the DFP procedure, after the kth step from  $x_k$  to  $x_{k+1}$ , a new estimate  $H_{k+1}$  to the inverse Hessian is sought, which is to replace the current estimate  $H_k$ . This new estimate is required to satisfy the quasi-Newton condition (also known as the DFP condition):

where  $\bar{y}_k$  is defined as  $(\bar{g}_{k+1} - \bar{g}_k)$ .

Where does this constraint come from? Basically, it is an *identity* which holds for *quadratic* functions. At the beginning of the kth iteration, we have a quadratic approximation to f(x), say:

(2.2) 
$$Q_k(x) = a_k + b_k^T x + \frac{1}{2} x^T G x,$$

(where the superscript T denotes the vector transpose) and, during this iteration, we make a step from the point  $x_k$  to a point  $x_{k+1}$ . At these two points, we have evaluated the exact gradient vectors:

(2.3) 
$$\bar{g}_k \equiv \nabla f(x_k); \quad \bar{g}_{k+1} \equiv \nabla f(x_{k+1}).$$

A new, improved quadratic approximation  $Q_{k+1}(x)$  is now forced to fit f(x) at these points, in the sense that the gradients calculated from  $Q_{k+1}(x)$  match the exact ones:

(2.4a) 
$$g_{k+1}(x_k) = b_{k+1} + G_{k+1}x_k = \bar{g}_k,$$

(2.4b) 
$$g_{k+1}(x_{k+1}) = b_{k+1} + G_{k+1}x_{k+1} = \bar{g}_{k+1}.$$

It follows that the new  $G_{k+1}$  satisfies the condition:

(2.5) 
$$\bar{g}_{k+1} - \bar{g}_k = G_{k+1}(x_{k+1} - x_k) \equiv G_{k+1}\sigma_k$$

which is equivalent to Eq. (2.1).

The method used in [7] to derive correction formulas was briefly as follows: The correction to  $H_k$  was written as:

$$(2.6) H_{k+1} = H_k + E_k$$

and a quadratic norm of  $E_k$  was minimized subject to (2.1). (This amounts to a constraint on  $E_k$ .) In addition, it was required that  $E_k$  be symmetric so as to preserve the symmetry of  $H_{k+1}$ , given that of  $H_k$ . This amounts to another (linear) constraint on  $E_k$ . This constrained variational problem was solved, leading to a class of correction formulas. These formulas resemble the DFP formula, and it was, in fact, shown by D. Goldfarb [13] that the variationally derived class contains the DFP formula.

3. Constraints in the Derivativeless Case. We now have the task of trans-

lating the variational procedure to the case when there is no independently calculated gradient. The first thing we must do is to find an appropriate constraint corresponding to the QN condition.

Clearly, the new condition cannot contain  $\bar{g}$  explicitly, since  $\bar{g}$  cannot be independently computed. Hence, the only admissible ingredients are the values of f at various points.

As in all treatments of quasi-Newton methods, we assume f(x) to be approximated by a quadratic function (as indicated previously). The approximation for f is Q, given in (2.2). If we replace b in favor of g, we obtain:

(3.1) 
$$Q = a + g^T x - \frac{1}{2} x^T G x.$$

This form for Q(x) has turned out, in practice, to be more convenient (less subject to rounding error) than that in (2.2), but it must be remembered that g depends on x.

Let us now assume that we are at some point  $x_0$  and do a line search along some s (with length parameter  $\alpha$ ) for the minimum of f. For any x on this line ( $x = x_0 + \alpha s$ ), we have for the estimate g, based on Q:

(3.2) 
$$g = b + Gx = b + G(x_0 + \alpha s)$$
  
=  $(b + Gx_0) + G(\alpha s) = g_0 + \alpha Gs.$ 

Correspondingly, for Q:

(3.3) 
$$Q = a + (g_0 + \alpha Gs)^T (x_0 + \alpha s) - \frac{1}{2} (x_0 + \alpha s)^T G(x_0 + \alpha s)$$
$$= (a + g_0^T x_0 - \frac{1}{2} x_0^T Gx_0) + (g_0^T s) \alpha + \frac{1}{2} (s^T Gs) \alpha^2.$$

At some value  $\alpha_1$ , we find the minimum value  $f_1$ . The corresponding x value is  $x_1 (=x_0 + \alpha_1 s)$ .

The spirit of the QN condition in the DFP case is to require that the estimated set of "parameters"  $\{H_{ik}\}$  be such as to make the quadratic representation Q "fit" the independently computed gradients. What corresponds in the present case is to require the "parameters"  $g_0$  and G to be such as to make the function Q(x) "fit" the independently computed values of f. Thus, we shall require for our next estimates,  $g_0^*$  and  $G^*$ , say:

(3.4a) 
$$Q_0 \equiv Q(0) = a + g_0^{*T} x_0 - \frac{1}{2} x_0^T G^* x_0 = f_0,$$

$$(3.4b) Q_1 \equiv Q(\alpha_1) = f_1.$$

As in the DFP method, we eliminate what amounts to an additive constant (viz., a) by taking differences:

(3.5) 
$$\Delta f = \Delta Q = Q_1 - Q_0 = (g_0^* s) \alpha_1 + \frac{1}{2} (s^T G^* s) \alpha_1^2.$$

There is another independent constraint, based on the fact that f is a minimum at  $\alpha_1$ . Hence, the derivative of Q, with respect to  $\alpha$ , is forced to vanish at  $\alpha_1$ :

(3.6) 
$$\left(\frac{dQ}{d\alpha}\right)_{\alpha_1} = g_0^{*T}s + (s^TG^*s)\alpha_1 = s^T(g_0^* + \alpha_1G^*s) = s^Tg_1^* = 0.$$

Thus, we have *two* "QN conditions" at each step. Other combinations are possible, of course, such as fitting  $Q(\alpha)$  to f at three distinct points along s. (This would also lead to two conditions.)

JOHN GREENSTADT

For reasons which will be apparent later, it is not feasible to attempt to correct  $g_0$  and G after only one step. We therefore take more steps than one in each "correction cycle", and distinguish between a *minor step*, involving a line search along a single direction, and a *major step*, which will be a sequence of such minor steps.

In what follows, we shall suppress the major step index k, and concentrate on the set of minor steps which constitute a major step.

Starting from  $x_0$  (the starting point of a major step), the first minor step direction  $d_1$  is calculated by Newton's formula, using the current estimates  $g_0$  and G:

$$(3.7) d_1 = -G^{-1}g_0$$

and  $d_1$  is then normalized with respect to a positive-definite matrix L, to be chosen later. This gives the unit vector  $s_1$ , defined as follows:

(3.8) 
$$s_1 \equiv d_1/(d_1^T L d_1)^{1/2}.$$

Note that it is necessary to solve a simultaneous linear system for  $d_1$ , since  $G^{-1}$  will not be directly estimated, as in the DFP method. The reason for this is that G is involved in Eqs. (3.5) in such a way, that replacing it by  $H^{-1}$  would unavoidably lead to a nonlinear constraint on H, thus rendering the variational problem intractable.

After the line search along  $s_1$ , yielding  $\alpha_1$  and  $f_1$ , the direction of the next minor step may be generated by combining  $s_1$  with some other direction. A simple choice is one of the coordinate directions, say  $e_1$ . Then

$$(3.9) d_2 = e_1 + \rho_1 s_1$$

with  $\rho_1$  chosen so as to make  $d_2$  orthogonal to  $s_1$ , in the sense that  $d_2^T L s_1 = 0$ .  $d_2$  is then normalized to give  $s_2$ , and a line search is performed, yielding  $\alpha_2$  and  $f_2$ . Next, a new direction  $d_3$  is found by combining  $e_2$ ,  $s_1$  and  $s_2$  linearly, and requiring  $d_3$  to be orthogonal to  $s_1$  and  $s_2$  (with respect to L).  $d_3$  is then normalized, etc.

If it should happen that one of the coordinate directions is a linear combination of the already computed direction vectors, it is simply dropped. In all, a total of N minor steps are attempted. In what follows, the index i will be a label for the minor steps within a major step.

If we denote the *i*th minor step by  $\sigma_i$ , we have:

$$(3.10) x_i = x_{i-1} + \sigma_i.$$

 $\tau_i$  is next defined as the total displacement from  $x_0$  to  $x_i$ :

(3.11) 
$$\tau_i \equiv x_i - x_0 = \sum_{j=1}^i \sigma_j.$$

Then, based on (3.2), we will impose the condition:

(3.12) 
$$g_i^* = g_0^* + G^* \tau_i = g_{i-1}^* + G^* \sigma_i$$

and, corresponding to (3.6), we satisfy:

(3.13) 
$$\sigma_{i}^{T}g_{i}^{*} = \sigma_{i}^{T}(g_{0}^{*} + G^{*}\tau_{i}) = 0$$

for each  $\tau_i$ .

Corresponding to (3.5), we have:

(3.14)  

$$\Delta f_{i} = f_{i} - f_{i-1} = Q_{i} - Q_{i-1}$$

$$= g_{i}^{*T} (x_{i-1} + \sigma_{i}) - (g_{i}^{*} - G^{*} \sigma_{i})^{T} x_{i-1}$$

$$- \frac{1}{2} (x_{i-1} + \sigma_{i})^{T} G^{*} (x_{i-1} + \sigma_{i}) + \frac{1}{2} x_{i-1}^{T} G^{*} x_{i-1}$$

$$= g_{i}^{*T} \sigma_{i} - \frac{1}{2} \sigma_{i}^{T} G^{*} \sigma_{i} = -\frac{1}{2} \sigma_{i}^{T} G^{*} \sigma_{i},$$

the last equation resulting from (3.13).

In summary, our constraints are:

$$(3.15a) \qquad \qquad \Delta f_i + \frac{1}{2} \sigma_i^T G^* \sigma_i = 0,$$

(3.15b) 
$$\sigma_i^T g_0^* + \sigma_i^T G^* \tau_i = 0.$$

It is important to note that the only independently computed functional quantities here are the  $\{\Delta f_i\}$ .

We are now going to consider the major step as an independent cycle, and make the corrections to our old estimates,  $g_0$  and G, at the end of it. The corrections will be denoted by  $\gamma$  and  $\Gamma$ , so that the corrected values  $g_0^*$  and  $G^*$  will be:

(3.16a) 
$$g_0^* = g_0 + \gamma$$
,

$$G^* = G + \Gamma.$$

Then the constraints (3.15), considered to apply to the new estimates  $g_0^*$  and  $G^*$ , are translated into constraints on  $\gamma$  and  $\Gamma$  as follows:

(3.17a) 
$$\frac{1}{2}\sigma_i^T\Gamma\sigma_i = -\{\Delta f_i + \frac{1}{2}\sigma_i^TG\sigma_i\} \equiv \rho_i$$

(3.17b) 
$$\sigma_i^T \gamma + \sigma_i^T \Gamma \tau_i = -\{\sigma_i^T g_0 + \sigma_i^T G \tau_i\} \equiv \epsilon_i.$$

Now, there are N parameters in  $g_0$  and  $\frac{1}{2}N(N + 1)$  in G to be estimated. But in each major step, we have at most 2N constraints. Hence, when N > 1, there are fewer constraints than parameters; so that one major step does not determine all the parameters. Since each major step is treated independently of the others, any method based on these constraints will not necessarily be an "N-step" method. In fact, the formulas to be derived need not necessarily generate the exact G, even for quadratic functions. This is not to say, however, that it is impossible to construct "N-step" formulas (by other means).

4. The Variational Procedure for the Derivativeless Case. We now have the problem of setting up a functional to minimize, which somehow embodies, the norms of  $\gamma$  and of  $\Gamma$ . The most obvious norms to choose, which are quadratic, are:

$$(4.1a) ||\gamma||^2 \equiv \gamma^T V\gamma,$$

(4.1b) 
$$||\Gamma||^2 \equiv \operatorname{Tr}(W\Gamma W\Gamma^T),$$

where V and W are positive-definite matrices of some sort.

A difficulty arises in somehow combining these norms in a natural manner. One wishes to have a quadratic function of the elements of  $\gamma$  and  $\Gamma$  which is also positivedefinite. These two quantities are not really comparable, since it is easy to construct functions for which they have arbitrary values. The obvious device of simply adding them leads to the problem of insuring that their "units" are consistent. This might be accomplished, for example, by taking  $W = G^{-1}$  and  $V = ||g_0||^{-2}I$ , where I is the unit matrix.

The most practical form, which was found after some trials, was the most obvious one, viz., a simple sum:

(4.2) 
$$\Phi_0 = \frac{1}{2} \gamma^T V \gamma + \frac{1}{2} \operatorname{Tr}(W \Gamma W \Gamma^T),$$

and a large number of numerical trials, wherein various forms of V and W were chosen, seemed to indicate that the choices V = I,  $W = \nu I$  (where  $\nu$  is some arbitrary number) worked best in practice. However, we shall defer this specialization to a later section, but leave V and W arbitrary so as to show the general form of the corrections.

Incorporating the constraints (3.17) into the functional via the Lagrange multipliers  $\{\eta_i\}$  and  $\{\theta_i\}$  gives:

(4.3)  
$$\Phi = \Phi_0 - \sum_i \eta_i (\frac{1}{2} \sigma_i^T \Gamma \sigma_i - \rho_i) - \sum_i \theta_i (\sigma_i^T \gamma + \sigma_i^T \Gamma \tau_i - \epsilon_i)$$

We should add to this the additional constraint  $\Gamma^{T} = \Gamma$ , but will dispense with doing this explicitly, and simply indicate the change in the formula for  $\Gamma$ , necessary to include this requirement.

The necessary conditions for a stationary  $\Phi$  are obtained by differentiating, as follows:

(4.4a) 
$$\frac{\partial \Phi}{\partial \gamma} = V\gamma - \sum_i \theta_i \sigma_i = 0,$$

(4.4b) 
$$\frac{\partial \Phi}{\partial \Gamma} = W \Gamma W - \sum_{i} \eta_{i} \cdot \frac{1}{2} \sigma_{i} \sigma_{i}^{T} - \frac{1}{2} \sum_{i} \theta_{i} (\sigma_{i} \tau_{i}^{T} + \tau_{i} \sigma_{i}^{T}) = 0.$$

(The symmetrizing of the  $\sigma_i \tau_i^T$  term is a result of taking account of the symmetry condition on  $\Gamma$ .)

If we define  $\Lambda \equiv V^{-1}$ ,  $M \equiv W^{-1}$ , we have:

(4.5a) 
$$\gamma = \Lambda \sum \theta_i \sigma_i$$

(4.5b) 
$$\Gamma = \frac{1}{2} M \{ \sum \eta_i \sigma_i \sigma_i^T + \sum \theta_i (\sigma_i \tau_i^T + \tau_i \sigma_i^T) \} M.$$

We now solve for the Lagrange multipliers  $\{\eta_i\}$  and  $\{\theta_i\}$  by applying the constraints to  $\gamma$  and  $\Gamma$ . The resulting equations are rather complicated, but they reduce to the following (in matrix form):

(4.6) 
$$A\theta + B\eta = \epsilon, \quad B^T\theta + C\eta = \rho,$$

where

(4.7) 
$$\epsilon \equiv {\epsilon_i}, \quad \rho \equiv {\rho_i},$$

(4.8a) 
$$A_{ij} \equiv \lambda_{ij} + \frac{1}{2} \{ \mu_{ij}^{(1)} \mu_{ij}^{(3)} + \mu_{ij}^{(2)} \mu_{ji}^{(2)} \},$$

(4.8b) 
$$B_{ij} \equiv \frac{1}{2} \mu_{ij}^{(2)} \mu_{ij}^{(3)},$$

(4.8c) 
$$C_{ij} \equiv \frac{1}{4} \mu_{ij}^{(3)} \mu_{ij}^{(3)}$$

and

(4.9a) 
$$\lambda_{ij} \equiv \sigma_i^T \Lambda \sigma_j,$$

$$(4.9b) \qquad \qquad \mu_{ij}^{(1)} \equiv \tau_i^T M \tau_{ij}$$

(4.9c) 
$$\mu_{ij}^{(2)} \equiv \tau_i^T M \sigma_j,$$

(4.9d) 
$$\mu_{ii}^{(3)} \equiv \sigma_i^T M \sigma_{ii}$$

i and j run from 1 to N and are not summed in (4.8).

If M and  $\Lambda$  are now chosen to be proportional to L, we gain a great simplification in the formulas for  $\gamma$  and  $\Gamma$ . We set (as suggested previously):

(4.10) 
$$W = v V; \text{ or } M = \frac{1}{v} \Lambda$$

and, in addition:

(4.11) 
$$\Lambda = L, \text{ so that } M = \frac{1}{\nu} L.$$

We then have, since  $\{s_i\}$  is now an orthonormal set with respect to L:

(4.12a) 
$$\lambda_{ij} = \sigma_i^T \Lambda \sigma_j = |\sigma_i| |\sigma_j| s_i^T L s_j = \sigma_i^2 \delta_{ij}$$

and, similarly:

(4.12b) 
$$\mu_{ij}^{(3)} = \sigma_i^T M \sigma_j = \frac{1}{\nu} \sigma_i^T L \sigma_j = \frac{1}{\nu} \sigma_i^2 \delta_{ij},$$

so that  $\{\lambda_{ij}\}$  and  $\{\mu_{ij}^{(3)}\}$  are diagonal. Since, from Eq. (3.11),  $\tau_i = \sum_{p=1}^i \sigma_p$ , we have:

(4.13)  
$$\mu_{ij}^{(2)} = \frac{1}{\nu} \sum_{p=1}^{i} \sigma_p^T L \sigma_j = \frac{1}{\nu} \sum_{p=1}^{i} \sigma_p^2 \delta_{pj}$$
$$= \sigma_j^2 / \nu, \quad \text{if } i \ge j,$$
$$= 0, \qquad \text{if } i < j,$$

so that  $\{\mu_{ij}^{(2)}\}$  is a lower triangular matrix.

Bearing in mind that the products in Eq. (4.8) are not matrix products, but element-by-element products, we see that:

- 1.  $\{\mu_{ij}^{(1)}\mu_{ij}^{(3)}\}$  is diagonal because  $\{\mu_{ij}^{(3)}\}$  is;
- 2.  $\{\mu_{ij}^{(2)}\mu_{ji}^{(2)}\}$  is diagonal because  $\{\mu_{ij}^{(2)}\}$  is triangular;
- 3.  $\{\mu_{ij}^{(2)}, \mu_{ij}^{(3)}\}\$  is diagonal because  $\{\mu_{ij}^{(3)}\}\$  is.

Hence,  $A_{ij}$ ,  $B_{ij}$  and  $C_{ij}$  all form diagonal matrices, and have the values:

(4.14a) 
$$A_{ij} = \left\{\sigma_i^2 + \frac{1}{2\nu^2} \left(\tau_i^2 \sigma_i^2 + \sigma_i^4\right)\right\} \delta_{ij},$$

(4.14b) 
$$B_{ij} = \frac{1}{2\nu^2} \sigma_i^4 \delta_{ij},$$

(4.14c) 
$$C_{ij} = \frac{1}{4\nu^2} \sigma_i^4 \delta_{ij},$$

where

(4.15) 
$$\tau_i^2 = \sum_{p=1}^i \sigma_p^2,$$

all of which follows from the orthonormality of  $\{s_i\}$  with respect to L.

The solution of Eq. (4.6) has the form:

(4.16a) 
$$\theta = (A - BC^{-1}B^{T})^{-1}(\epsilon - BC^{-1}\rho),$$

(4.16b) 
$$\eta = C^{-1}(\rho - B^T \theta)$$

and these expressions may be easily evaluated because all the matrices are diagonal. The result is (by components):

(4.17a) 
$$\theta_i = \frac{2\nu^2(\epsilon_i - 2\rho_i)}{\sigma_i^2(2\nu^2 + \tau_i^2 - \sigma_i^2)},$$

(4.17b) 
$$\eta_i = \frac{4\nu^2}{\sigma_i^4} \rho_i - 2\theta_i,$$

so that the evaluation of  $\gamma$  and  $\Gamma$  does not really involve any matrix inversions.

The algorithm now runs as follows:

1. Assume G = I, and estimate  $g_0$  at the starting point by first differences. (See explanation in Section 5.)

2. To start a major step, compute a direction  $s_1$  from Eqs. (3.7), (3.8).

3. Do a line search for a minimum of f along s (for each minor step).

4. Save  $\sigma$ ,  $\tau$ ,  $\rho$  and  $\epsilon$  as defined in Section 3. If a total of N independent directions have been generated, skip to step 6.

5. Form a new direction from the previous step directions plus a new linearly independent direction, and orthonormalize. Go to step 3.

6. Compute  $\theta$  and  $\eta$  from Eqs. (4.17).

7. Compute  $\gamma$  and  $\Gamma$  from (4.5).

8. Correct  $g_0$  and G (Eq. (3.16)) to form  $g_0^*$  and  $G^*$ .

9. Translate  $g_0^*$  using  $g_0^{**} = g_0^* + G^* \tau_N$  (referring to Eq. (3.12), since the new  $x_0^*$  is  $x_0 + \tau_N$ ).

This completes a major step.

10. Test for termination  $(||g_0^{**}|| < \text{threshold, say})$ . Otherwise, go back to step 2. There are the usual complications in the program for this algorithm, mostly as a result of rounding error. These have not been described here.

5. Computational Experience. This method was programmed in the APL language for the IBM 360 computer and a good many trials were run on a few test functions. There was a good deal of tinkering necessary to get the method to converge reliably and reasonably efficiently, but the most effective choice of various arbitrary quantities turned out to be one of the simplest.

The worst difficulty with this method is that the successive estimates of G are not necessarily positive-definite. This precludes setting L = G (hence  $\Lambda = G$  and  $M = G/\nu$ ) since minimizing a quadratic form with an indefinite metric can (and did !) yield very large, unstable corrections  $\gamma$  or  $\Gamma$ . The choice L = I turned out to be the most stable (and the simplest) choice, and almost always led to the fastest convergence.

The best choice of  $\nu$  turned out to be 0! Of course, one cannot simply set  $\nu = 0$ , and evaluate  $\gamma$  and  $\Gamma$ , since Eqs. (4.6) become singular for  $\nu = 0$ . It is possible, of course, to find the limiting solution as  $\nu \to 0$ , and this is described in Appendix A.

In many instances, the correction computed in this way caused G to become indefinite. This is easily detectable in those cases when a diagonal element becomes negative. This was cleared up in most instances by letting  $\nu \to \infty$ , instead of  $\nu \to 0$ . (The former case is analyzed in Appendix A.) When this device did not help, the indefiniteness was allowed to remain, and the next major step was begun. Near the point of convergence, this pathological effect nearly always disappeared; however, it did have the effect of slowing down convergence.

As will be seen from the printouts of some of the examples shown, the convergence does seem to be superlinear in many cases. This has not been proven and may not even always be true.

There is certainly no assurance that a variational derivation will yield formulas having the most desirable properties. It is likely that a deeper theoretical analysis of this type of QN method will yield better procedures with better properties (such as positive-definite G's).

As in the DFP method, the unit matrix was taken as a starting value for G. For a starting value of  $g_0$ , there is no "natural" vector, although, in principle, it is possible to start with any vector. When this was done (for example, by taking  $g_0 = (10000 \cdots)$  or  $g_0 = (111 \cdots)$ , the method converged, but often with great difficulty. Ultimately, a rough estimate of  $g_0$  was computed at the outset (by simple forward differences), and this stabilized matters quite considerably.

6. Numerical Examples. Tables 1-3 following are printouts generated at a terminal by the APL program. The entries are as follows:

NSTEP	The	maior	step	number.
		major	o vo p	mannoer,

- P The number of minor steps in the major step; in these tables, P = Nin all cases, except when some minor steps are too small. (The formulas for  $\theta$  and  $\eta$  remain the same, except that N is replaced by P.)
- NFUNC The total number of evaluations of f after each major step.
  - F The value of f(x).
  - X The position vector.

In these printouts,  $g_0$  is denoted by GZ and G is denoted by GG. When G is found to be indefinite, the notation: IG (indefinite G) with the major step number is printed. The value of  $\nu$  is then changed from 0 to  $\infty$ . When this still gives a detectably indefinite G, the same notation is printed again. The entire process was regarded as having converged when  $||g_0|| < 10^{-5}$ , or, failing this, that no minor step  $> 10^{-7}$  was possible. If the size of the major step falls below  $10^{-6}$ , the notation "SPF" is printed, and the iteration terminated.

The functions tested were as follows: (The starting values in each case are listed on the first line with NSTEP = 0.)

(a) Quadratic Function 1.

$$f = x_1^2 + 100(x_2 - 1)^2 + (x_3 - 2)^2$$

whose Hessian is equal to:

$$G_1 = \begin{bmatrix} 2 & 0 & 0 \\ 0 & 200 & 0 \\ 0 & 0 & 2 \end{bmatrix}$$

and minimum at (0, 1, 2). Various starting values were used.

(b) Quadratic Function 2.

$$f = (x_1 + x_2 - 2)^2 + 10^4 (x_1 - x_2)^2$$

with Hessian:

$$G_2 = \begin{bmatrix} 20002 & -19998 \\ -19998 & 20002 \end{bmatrix}$$

and minimum at (1, 1).

(c) Quadratic Function 3.

$$f = (x_1 + 2x_2 + 3x_3)^2 + 100(x_2 - 1)^2 + (x_3 - 2)^2,$$
$$G_3 = \begin{bmatrix} 2 & 4 & 6 \\ 4 & 208 & 12 \\ 6 & 12 & 20 \end{bmatrix}$$

and minimum at (-8, 1, 2).

(d) Rosenbrock's Function [8].

$$f = 100(x_2 - x_1^2)^2 + (1 - x_1)^2,$$
  

$$G_{\text{Ros}} = \begin{bmatrix} 802 & -400 \\ -400 & 200 \end{bmatrix} \text{ at } (1, 1).$$

(e) Beales's Function [9].

$$f = \sum_{i=1}^{3} [c_i - x_1(1 - x_2^i)]^2; \quad \{c_i\} = \{1.5, 2.25, 2.625\}$$

(Hessian not computed independently).

(f) Powell's Function No. 1 [10].

$$f = (x_1 + 10x_2)^2 + 5(x_3 - x_4)^2 + (x_2 - 2x_3)^4 + 10(x_1 - x_4)^4$$

(Hessian not computed independently).

(g) Powell's Function No. 2 [11].

$$f = \left[1 + (x_1 - x_2)^2\right]^{-1} + \sin(\frac{1}{2}\pi x_2 x_3) + \exp\left\{-\left(\frac{x_1 + x_3}{x_2} - 2\right)^2\right\}$$

1

(Hessian not computed independently).

(h) Cube [12].

$$f = 100(x_2 - x_1^3)^2 + (1 - x_1)^2,$$
  

$$G_{\text{CUBE}} = \begin{bmatrix} 1802 & -600 \\ -600 & 200 \end{bmatrix} \text{ at } (1, 1).$$

(i) Random Trigonometric Function [2].

$$f = \sum_{i=1}^{N} \left\{ E_i - \sum_{j=1}^{N} (A_{ij} \sin x_j + B_{ij} \cos x_j) \right\}^2$$

(with  $A_{ij}$ ,  $B_{ij}$ , and  $E_i$  randomly generated).

 $G_{RT}$  is variable and the solution is "XNULL", which is precomputed.

(j) Helical Valley [2].

$$f = 100[(x_3 - 10\theta)^2 + (r - 1)^2] + x_3^2$$

with

$$\theta = \tan^{-1}(x_2/x_1); r = (x_1^2 + x_2^2)^{1/2}$$

Solution: (1, 0, 0).

(k) Wood's Function [15].

$$f = 100(x_2 - x_1^2)^2 + (1 - x_1)^2 + 90(x_4 - x_3^2)^2 + (1 - x_3)^2 + 10.1\{(x_2 - 1)^2 + (x_4 - 1)^2\} + 19.8(x_2 - 1)(x_4 - 1).$$

Solution: (1, 1, 1, 1).

It will be seen that various interesting (some good and some bad) things occur in these problems:

(1) The convergence near the solution is often clearly superlinear (even quadratic at times), but breaks down for functions which do not have a quadratic minimum (e.g., Powell 1).

(2) When G at the solution is singular, there is a good deal of difficulty with indefinite intermediate G's, which slows the convergence drastically.

(3) This method is not as speedy as several others (Simplex, Powell's, Rosenbrock's) but compares well in some cases.

(4) The successive estimates of G have been printed for Quadratic Function No. 3; evidently, a good value is generated very soon, which explains the quite rapid convergence in the quadratic cases. (A similar study of what happens to  $g_0$  has not been made.)

(5) When P, the number of minor steps per major step is restricted to be  $\langle N$ , the convergence is slowed considerably. (These cases are not shown.) When P = 1, the correction to  $g_0$  tends to make it vanish altogether, thus providing no direction for the next Newton step. (This was the reason for introducing additional minor steps in the first place.)

In Table 4 is shown a comparison with other methods for those test functions for which information is available. The starting points for all comparison functions are the "standard" ones, i.e., those used most in the literature.

The entries in Table 4 are as follows:

QNWD stands for "Quasi-Newton Without Derivatives"

- H-J stands for "Hooke and Jeeves"
- Ros stands for "Rosenbrock"
- SPLX stands for "Simplex"
- Pow stands for "Powell"
- Stew stands for "Stewart".

For each case, the number of function evaluations needed to get the function

```
TABLE 1
```

QUADR. FUNC 2 F 1.0001E0W 9.3134E<sup>-</sup>01 2.3272E<sup>-</sup>09 W.9420E<sup>-</sup>17 X+ 1.0006F02 0.0000600 1.0096F00 9.9937F01 1.00006E0 1.0000600 1.0000600 1.0000600 NSTEP P NPUNC 
 STLP
 P
 NFUNC

 Q
 0
 3

 1
 2
 15

 2
 1
 25

 3
 1
 35

 CONVERGED
 3
 0

 3
 0
 45
 GZ 45 #.9414E<sup>-</sup>17 1.0000E00 1.0000E00 -0.47993 0.47984 GG 20002 <sup>-</sup>19998 -19998 20002 QUADR. FUNC. 3 
 #STEP
 N.FUNC
 F
 X+

 0
 0
 2.0100E02
 3.0000E00
 2.0000E00
 1.0000E00

 1
 3
 15
 1.2464E01
 -4.4440E00
 1.1904E00
 1.6976E00

 NJF(91)
 -</ ZW[91] 2.0573 3.0072 6.0731 3.0072 200 12.065 6.0731 12.065 19.942 2 3 25 1.43968<sup>-</sup>02 7.6544800 9.97218<sup>-</sup>01 1.9123800 1.9107 4.4548 5.8024 4.4548 208.09 11.949 5.8024 11.949 20.003 3 34 1.90525<sup>-</sup>05 <sup>-7</sup>.9869500 1.00005<sup>-</sup>00 1.9956500 з 2.0031 4.0165 6.005# 4.0165 20# 12.056 6.005# 12.056 19.997 2 47 1.1494**5**11 #.0000**5**00 \* 1.0000**5**00 2.0000500 ٠ 1.9942 4.0165 5.9977 4.0165 208 12.056 5.9977 12.056 20.002 5 1 60 3.40658 17 4.0000200 1.0000200 2.0000200 \*SPF computer CONVERGED CONVERGED 5 1 60 3.4065E<sup>-</sup>17 8.0000E00 1.0000E00 2.0000E00 CZ <sup>-</sup>4.343WE<sup>-</sup>9 <sup>-</sup>2.1134E<sup>-</sup>7 8.6236E<sup>-</sup>10 GG 1.9985 4.0165 5.997 4.0165 208 12.056 5.9976 12.056 20.002 5.9976 

 BSTEP
 P
 FUNC
 P
 X+

 0
 0
 4
 85440E03
 1.0000E01
 1.0000E01
 1.0000E01

 1
 3
 15
 2.1194E02
 W.2552E00
 2.4072E00
 W.3553E500

 2
 3
 2.5
 4.3574E
 0.3
 7.4742E00
 1.0014E00
 1.9554E00

 3
 3
 4.6.0971E
 0.4
 7.9214E00
 9.9995E
 01
 1.9554E00

 4
 3
 3.0676E
 1
 W.0000E00
 1.0000E
 0.0000E00
 2.0000E00
 5
 1
 56
 1.0791E
 1.4
 W.0000E00
 1.0000E
 0.000E
 2.0000E00

 5
 1
 56
 1.0791E
 W.0000E00
 1.0000E
 0.000E
 2.0000E00

 5
 1
 56
 1.0791E
 W.0000E00
 1.0000E
 0.000E
 0.000E
 0.000E

 5
 1
 56
 1.0791E
 W.0000E00
 1.0000E
 0.000E
 0.000E
 0.000E
 0.000E

 6
 1
 56
 1.0791E
 W.0000E00
 1.0000E
 0.000E
 0.000E
 0.000E GZ 7.2411E 8 1.6262E 6 1.9152E 8 GG 1.9978 3.8979 5.9977 3.8979 208 11.659 5.9977 11.659 20.002 QUADR. FUNC. 3 NSTEP P NFUNC 
 O SEP
 P
 NFUNC

 0
 0
 4

 1
 3
 14

 2
 3
 3

 3
 3
 40

 4
 2
 54

 CONVERCED
 66
 F X+ 1.0104E04 1.0000E02 0.0000E00 0.0000E00 2.8439E03 4.9717E01 4.2341E<sup>-</sup>01 9.3294E<sup>-</sup>01 8.3338E<sup>-</sup>03 <sup>-</sup>8.0426E00 1.0002E00 1.9841E00 1.3427E<sup>-</sup>06 <sup>-</sup>8.0037E00 1.0000E00 2.0000E00 2.0308E<sup>-</sup>13 <sup>-</sup>8.0000E00 1.0000E00 2.0000E00 0 66 1.3073E<sup>-</sup>13 <sup>-</sup>8.0000E00 1.0000E00 2.0000E00 2.7719E 7 1.8613E 5 6.874E H CC 1.9998 3.9947 5.999 3.9947 208 12.011 5.9997 12.011 20 5.9997

# TABLE 1 (continued)

QUADR. FUNC.	1			
NSTEP P NFUNC	F	X+	0. 0000F00	1
1 3 14	1.1000A02 8.7319E 01	2.6728E 01	1.0891200	1.9109200
2 2 31 3 2 49	1.4443E-10 2.55805-12	9.4633E 06 1.5982E 06	1.0000E 00 1.0000E00	2.0000E00 2.0000E00
+SPF 3 Converged		_		
3249 GZ	2.55802 12	1.5982 <i>E</i> 06	1.0000500	2.0000800
1.80115 0.00	065156 5.7706 <i>2</i>	10		
GG				
2.000050 1	4511E 4 1.814	0E 4		
1.81405 4 1	.36558 4 2.000	OE O		
QUADR. FUNC. 1	ч			
NSTEP P NFUNC	F 8.3440E03	X+ -1.0000E01	1.0000E01	-1.0000 <i>E</i> 01
1 3 15	6.8539200	2.86605 01	1.2579200	1.6561E00 2.0000E00
3 1 39	4.98392-14	-2.8514E-10	1.0000800	2.0000500
CONVERGED				
3 1 39 GZ	4.98392 14	2.85148 10	1.0000200	2.0000200
-4.3749E 6 0.00	039373 5.2498	8-6		
GG				
2.0000E0 1. 1.3838E 5 2.	383855 7.044 000052 1.603	48 <sup>-7</sup> 58 <sup>-</sup> 5		
7.044457 1.	6035E 5 2.000	0 E 0		
QUADR. PUNC.				
NSTEP P NEUNC	,	7+		
0 0 4	1.0104504	1.0000E02	0.0000 <u>5</u> 00	0.0000E00 7 HL7LE 02
2 3 29	2.77465 06	1.19075 03	9.99995-01	1.9988200
CONVERGED	1.68892 15	2.31142 09	1.0000200	2.0000200
3 1 43 GZ	1.68892-15	2.31145-09	1.0000500	2.0000200
3 1 43 <i>GZ</i> 7.8167 <i>E</i> <sup>-</sup> 10 2.6	1.6889E <sup>-</sup> 15 501E <sup>-</sup> 7 - 5.2693E	2.3114 <i>6</i> -09	1.0000500	2.0000200
3 1 43 GZ 7.8167E <sup>-</sup> 10 2.6 GG	1.6889E <sup>-</sup> 15 501E <sup>-</sup> 7 <sup>-</sup> 5.2693E	2.3114 <i>6</i> 09	1.0000500	2.0000200
3 1 43 GZ 7.0167E <sup>-</sup> 10 2.61 GG 1.9990E0 2 2.9797E <sup>-</sup> 5 2	1.6889E <sup>-</sup> 15 501E <sup>-</sup> 7 <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 2.209 .0000E2 <sup>-</sup> 6.680	2.31148-09 	1.0000 <i>E</i> 00	2.0000500
3 1 43 <i>GZ</i> 7.8167 <i>E</i> <sup>-</sup> 10 2.6 <i>GG</i> 1.999% <i>E</i> 0 2 2.9797 <i>E</i> <sup>-</sup> 5 2 2.2098 <i>E</i> <sup>-</sup> 4 6	1.68892 <sup>-</sup> 15 5012 <sup>-</sup> 7 <sup>-</sup> 5.26932 .97972 <sup>-</sup> 5 2.209 .00002 <sup>-</sup> 6.680 .68002 <sup>-</sup> 4 2.000	2.31148-09 	1.0000 <i>E</i> 00	2.0000 <i>E</i> 00
3 1 43 7.8167E <sup>-</sup> 10 2.69 <i>GG</i> 1.999%E0 2 2.9797E <sup>-</sup> 5 2 2.209%E <sup>-</sup> 4 <sup>-</sup> 6	1.6889E <sup>-</sup> 15 501E <sup>-</sup> 7 <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 2.209 .0000E <sup>2</sup> <sup>-</sup> 6.680 .6800E <sup>-</sup> 4 2.000	2.31145-09 	1.0000800	2.0000200
3 1 43 CZ 7.8167E <sup>-</sup> 10 2.63 GG 1.9998E0 2 2.9797E <sup>-</sup> 5 2 2.2098E <sup>-</sup> 4 - 6 QUADR. PUNC	1.6889E <sup>-</sup> 15 501E <sup>-</sup> 7 <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 2.209 .0000E <sup>2</sup> <sup>-</sup> 6.680 .6800E <sup>-</sup> 2.000	2.31145-09 	1.0000800	2.0000200
3 1 43 CZ 7.8167E <sup>-</sup> 10 2.63 GG 1.9998E0 2 2.9797E <sup>-</sup> 5 2 2.2098E <sup>-</sup> 4 -6 QUADR. PUNC MSTEP P NFUNC	1.6889E <sup>-</sup> 15 501E <sup>-</sup> 7 <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 2.209 .0000E <sup>2</sup> <sup>-</sup> 6.880 .6800E <sup>-</sup> , 2.000 2 P	2.3114E <sup>-</sup> 09 	1.0000800	2.0000200
3 1 43 CZ 7.8167E <sup>-</sup> 10 2.61 GG 1.999WE0 2 2.3797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 4 <sup>-</sup> 6 QUADR. FUNC NSTEP P NFUNC 0 0 3 1 2 9	1.6489E <sup>-</sup> 15 501E <sup>-</sup> 7 <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 2.209 .0000E <sup>2</sup> <sup>-</sup> 6.640 .6400E <sup>-</sup> 2.000 2 <i>P</i> 3.22405E02 3.2367E02	2.3114509 -10 -10 -10 -10 -10 -10 -10 -0000501 -9.9944500	1.0000E00 1.0001E01 9.935WE00	2.0000200
3 1 43 CZ 7.8167E <sup>-</sup> 10 2.61 GG 1.9998E0 2 2.9797E <sup>-</sup> 5 2 2.2098E <sup>-</sup> 4 <sup>-</sup> 6 QUADR. FUNC NSTEP P NFUNC 0 0 3 1 2 9 2 2 2 3 3 1 43	1.6489E <sup>-</sup> 15 501E <sup>-</sup> 7 <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 2.209 .0000E2 <sup>-</sup> 6.640 .6400E <sup>-</sup> 4 2.000 2 <i>P</i> 3.22405E02 3.2367502 1.3747500 6.9562E <sup>-</sup> 14	2.3114509 	1.0001201 9.995¥200 1.5862200	2.0000200
3 1 43 CZ 7.8167E <sup>-</sup> 10 2.61 GG 1.9998E0 2 2.9797E <sup>-</sup> 5 2 2.2098E <sup>-</sup> 4 -6 QUADR. PUNC NSTEP P NFUNC 0 0 3 1 2 9 2 2 2 3 3 1 43 CONVERGED 3 1 43	1.6489E <sup>-</sup> 15 501E <sup>-</sup> 7 <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 2.209 .0000E2 <sup>-</sup> 6.640 .6400E <sup>-</sup> 4 2.000 2 P 3.2405E02 3.2367E02 1.3747E00 6.9562E <sup>-</sup> 18 4.6627 <sup>-</sup> 18	2.3114509 -10 -10 -10 -10 -10 -10 -10 -10	1.0000200 1.0001201 9.95%200 1.5%2200 1.000200	2.0000200
3 1 43 CZ 7.8167E <sup>-</sup> 10 2.63 GG 1.999WE0 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 6 QUADR. PUNC NSTEP P NFUNC 0 0 3 1 2 9 2 2 2 3 3 1 43 CONVERGED 3 1 43	1.6449E <sup>-</sup> 15 501E <sup>-</sup> 7 <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 2.209 .0000E <sup>2</sup> <sup>-</sup> 6.640 .6400E <sup>-</sup> 4 2.000 2 2 3.2405E02 3.2367802 1.3747800 6.9562E <sup>-</sup> 14 2357 <sup>-</sup> 7	2.311+5°09 **10 **10 ***************************	1.0001201 9.95%200 1.5%2200 1.0000200	2.0000200
3 1 43 CZ 7.8167E <sup>-</sup> 10 2.63 GG 1.999WE0 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 6 QUADR. PUNC NSTEP P NFUNC 0 0 3 1 2 9 2 2 2 3 3 1 43 CONVERGED 3 1 43 CONVERGED CO	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>-</sup> 5.2693E .9797E <sup>-5</sup> 2.209 .0000E <sup>2</sup> <sup>-</sup> 6.640 .6400E <sup>-</sup> 4 2.000 2 <b>P</b> 3.2405E02 3.2367802 1.3747800 6.9562E <sup>-</sup> 14 235E <sup>-</sup> 7	2.311+5°09 **10 **50 **750 ***************************	1.0001201 9.9352200 1.5462200 1.0000200	2.0000200
3 1 43 CZ 7.8167E <sup>-</sup> 10 2.63 GG 1.999WE0 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 3.1 43 CONVERCED 3.1 43 GZ 9.6235E <sup>-</sup> 7 - 9.6 GC	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 <sup>2</sup> .209 .000052 <sup>-</sup> 6.640 .6400E <sup>-</sup> 4 <sup>2</sup> .000 2 <b>P</b> 3.2367E02 3.2367E02 1.3747E00 6.9562E <sup>-</sup> 14 235E <sup>-</sup> 7	2.311+F 09 	1.0001501 9.9954500 1.5462500 1.0000500	2.0000200
3 1 43 CZ CZ 7.8167E <sup>-</sup> 10 2.63 GC 1.999WE0 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 3.1 2	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 <sup>-</sup> 2.209 .000052 <sup>-</sup> 6.640 .6400E <sup>-</sup> 4 <sup>-</sup> 2.000 2 <sup>9</sup> <sup>9</sup> .2405E02 1.3747E00 6.9562E <sup>-</sup> 14 235E <sup>-7</sup> <sup>-</sup> 19998 20002	2.311+F 09 	1.0001501 9.9954500 1.5462500 1.0000500	2.0000200
3 1 43 CZ CZ 7.8167E <sup>-</sup> 10 2.61 GC 1.999WEO 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 CONVERCE 0 0 3 1 2 9 2 23 3 1 43 CONVERCED 3 1 43 9.6235E <sup>-</sup> 7 -9.6 GC 20002 -1999W	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>-</sup> 5.2693E .0000E <sup>2</sup> <sup>-</sup> 6.640 .6400E <sup>-</sup> 4 2.000 2 <b>P</b> 3.2405E02 3.2367E02 1.3747E00 6.9562E <sup>-</sup> 14 235E <sup>-7</sup> <sup>-</sup> 19994 20002	2.311+E <sup>-</sup> 09 	1.0000F00 1.0001F01 9.95%E00 1.5862E00 1.0000F00	2.0000200
3 1 43 CZ CZ 7.8167E <sup>-</sup> 10 2.63 GC 1.999WE0 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 CUADR. FUNC 0 3 1 2 9 2 2 3 3 1 43 CONVERCED 3 1 43 9.6235E <sup>-</sup> 7 -9.6 GC 20002 -1999W	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>-</sup> 5.2693E .0000E <sup>2</sup> <sup>-</sup> 6.640 .6400E <sup>-</sup> 4 2.000 2 <b>P</b> 3.2405E02 3.2367E02 1.3747E00 6.9562E <sup>-</sup> 14 235E <sup>-</sup> 7 <sup>-</sup> 19994 20002 2	2.311+E <sup>-</sup> 09 	1.0000F00 1.0001F01 9.95%F00 1.5862E00 1.0000F00	2.0000200
3 1 43 CZ CZ 7.8167E <sup>-</sup> 10 2.63 GC 1.999WE0 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 CUADR. FUNC 0 3 1 2 9 2 2 3 3 1 43 CONVERCED 3 2 2 3 3 1 43 9.6235E <sup>-</sup> 7 -9.6 GC 20002 -1999W QUADR. FUNC NSTEP P NFURC	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>5</sup> .2693E .0000E <sup>2</sup> <sup>6</sup> .640 .6400E <sup>-</sup> 2.000 2 <b>P</b> 3.2405E02 3.2367E02 1.3747E00 6.9562E <sup>-</sup> 14 235E <sup>-7</sup> <sup>1</sup> 19994 20002 2 <i>F</i>	2.311+E <sup>-</sup> 09 	1.0000F00 1.0001F01 9.9954F00 1.582E00 1.0000F00	2.0000200
3 1 43 CZ CZ 7.8167E <sup>-</sup> 10 2.63 GC 1.999WEO 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 - CUADR. FUNC 0 0 3 1 2 9 2 2 3 3 1 43 CONVERCED 3 2 2 3 3 1 43 9.6235E <sup>-</sup> 7 -9.6 GC 20002 -1999W QUADR. FUNC 0 0 3 1 2 12 NSTEP P NFUKC 0 0 3 1 2 12 2 2 2 2 2 2 3 1 43 CONVERCED 3 2 2 2 3 1 43 CONVERCED 3 1 43 CONVERCED 3 2 2 2 3 1 43 CONVERCED 3 2 2 2 3 1 43 CONVERCED 3 1 43 CONVERCED 3 2 2 2 3 1 43 CONVERCED 3 1 43 CONVERCED 3 2 2 2 3 1 43 CONVERCED 3 2 2 2 2 2 2 3 1 43 CONVERCED 3 2 2 2 CONVERCED 3 2 2 2 2 2 CONVERCED 3 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>-</sup> 5.2693E .9797E <sup>-</sup> 5 <sup>2</sup> .209 .0000E <sup>2</sup> <sup>-</sup> 6.640 .6400E <sup>-</sup> 4 <sup>2</sup> .000 2 <sup>9</sup> 3.2405E02 3.2367E02 1.3747E00 6.9562E <sup>-</sup> 14 235E <sup>-</sup> 7 <sup>-</sup> 19994 20002 2 <sup>7</sup> 4.0000E06 2.2491E <sup>-</sup> 06	2.311+E <sup>-</sup> 09 	1.0000F00 1.0001F01 9.95%F00 1.582E00 1.0000F00 1.0000F00	2.0000200
3 1 43 CZ CZ 7.8167E <sup>-</sup> 10 2.63 GC 1.999WEO 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 - CUADR. FUNC 0 0 3 1 2 9 2 2 3 3 1 43 CONVERCED 3 1 43 9.6235E <sup>-</sup> 7 -9.6 GC 20002 -1999W QUADR. FUNC 0 3 1 2 12 2 1 20 4 20 1 2 12 2 1 20 4 20 4 20 4 20 4 20 5	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>-5</sup> .2693E .9797E <sup>-5</sup> 2.209 .0000E <sup>2</sup> <sup>-6.64C</sup> .6400E <sup>-</sup> 4 2.000 2 <b>P</b> 3.2405E02 3.2405E02 3.2405E02 1.37+7500 6.9562E <sup>-</sup> 14 235E <sup>-7</sup> <sup>-19994</sup> 20002 2 <b>P</b> .0000E06 2.241E <sup>-</sup> 06 1.3917E <sup>-</sup> 13	2.311+E <sup>-</sup> 09 	1.0000F00 1.0001F01 9.95%F00 1.5852E00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00	2.0000200
3 1 43 CZ CZ 7.8167E <sup>-</sup> 10 2.63 GC 1.999WE0 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 - CUADR. FUNC 0 0 3 1 2 9 2 2 3 3 1 43 CONVERCED 3 1 43 9.6235E <sup>-</sup> 7 -9.6 GC 20002 -1999W QUADR. FUNC 0 3 1 2 2 2 3 1 43 CONVERCED 3 2 2 3 9.6235E <sup>-</sup> 7 -9.6 GC 20002 -1999W QUADR. FUNC 0 3 3 2 2 2 2 3 1 43 CONVERCED 2 2 3 2 2 0 3 7 2 2 1 20 2 2 3 1 2 1 20 2 2 3 3 1 3 1 43 CONVERCED 3 1 43 CONVERCED 2 1 20 CONVERCED 2 0 37 CONVERCED 2 0 37 CONVERCED CO	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>-5</sup> .2693E .9797E <sup>-5</sup> 2.209 .0000E <sup>2</sup> <sup>-6.64C</sup> .6400E <sup>-</sup> 2.000 2 <b>P</b> 3.2405E02 3.2405E02 3.247802 1.37+7500 6.9562E <sup>-</sup> 14 235E <sup>-7</sup> <sup>-19994</sup> 20002 2 <b>P</b> .0000E06 2.2491E <sup>-</sup> 06 1.3917E <sup>-</sup> 13 1.3970E <sup>-</sup> 15	2.311+E <sup>-</sup> 09 	1.0000F00 1.0001F01 9.95%F00 1.5852E00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00	2.0000200
3 1 43 CZ CZ 7.8167E <sup>-</sup> 10 2.63 GC 1.999WE0 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 - QUADR. FUNC 0 0 3 1 2 9 2 2 3 3 1 43 CONVERCED 3 1 43 CONVERCED 3 1 43 CONVERCED 3 1 43 CONVERCED 3 2 23 9.6235E <sup>-</sup> 7 -9.6 GC 20002 -1999W QUADR. FUNC 0 0 3 1 2 12 2 1 20 - 2 0 37 0 0 0 3 1 2 12 0 3 7 0 0 3 0 3 1 2 12 0 3 0 3 0 3 0 3 0 3 0 3 0 3 0 3	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>5</sup> .2693E .9797E <sup>-5</sup> <sup>2</sup> .209 .0000E <sup>2</sup> <sup>6</sup> .640 .6400E <sup>-</sup> 4 <sup>2</sup> .000 2 <sup>9</sup> .2405E02 1.37+7500 6.9562E <sup>-</sup> 14 6.9562E <sup>-</sup> 14 235E <sup>-7</sup> <sup>-19994</sup> 20002 2 <sup>9</sup> .2009E <sup>0</sup> 6 1.3917E <sup>-</sup> 13 1.3970E <sup>-</sup> 15 D09	2.311+E <sup>-</sup> 09 	1.0000F00 1.0001E01 9.95%F00 1.5852E00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00	2.0000200
3 1 43 CZ 7.8167E <sup>-</sup> 10 2.63 GC 1.999WE0 2 2.9797E <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 2 2.209WE <sup>-</sup> 5 - QUADR. FUNC 0 0 3 1 2 9 2 2 3 3 1 43 CONVERCED 3 1 43 9.6235E <sup>-</sup> 7 -9.6 GC 20002 -1999W QUADR. FUNC 0 0 3 1 2 12 2 1 20 1 2 12 2 1 20 1 2 12 2 0 37 0.00WERCED 2 0 37 0.00WERCED 2 0 37 0.00WERCED 2 0 37 0.00WERCED	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>-5</sup> .2693E .9797E <sup>-5</sup> 2.209 .0000E <sup>2</sup> <sup>-6.64C</sup> .6400E <sup>-4</sup> 2.000 2 <sup>9</sup> 3.2405E02 3.2405E02 3.2405E02 3.2405E02 3.2405E02 3.2405E02 3.2405E02 1.3747F00 6.9562E <sup>-144</sup> 6.9562E <sup>-144</sup> 235E <sup>-7</sup> <sup>-19994</sup> 20002 2 <sup>9</sup> <sup>10956</sup> 2.2491E <sup>-06</sup> 1.3917E <sup>-13</sup> 1.3970E <sup>-15</sup> DB9	2.311+E <sup>-</sup> 09 	1.0000F00 1.0001E01 9.95%F00 1.562E00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00	2.0000200
3 1 43 CZ CZ GG 1.999WE0 2 2.9797E5 2 2.209WE 5 2 2.209WE 5 6 QUADR. FUNC 0 0 3 1 2 9 2 2 3 3 1 43 CONVERCED 3 1 43 CONVERCED 3 1 43 9.6235E77 -9.6 GG 20002 1999W QUADR. FUNC 0 0 3 1 2 12 2 1 20 1 2 12 2 1 20 1 2 12 2 0 37 0.00WERCED 2 0 0 3 0 0 10 0 0 10 0 0 10 0 0 10 0 0 10 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	1.6449E <sup>-15</sup> 501E <sup>-7</sup> <sup>5</sup> .2693E .9797E <sup>-5</sup> <sup>2</sup> .209 .0000E <sup>2</sup> <sup>6</sup> .640 .6400E <sup>-</sup> 4 <sup>2</sup> .000 2 <sup>9</sup> <sup>2</sup>	2.311+E <sup>-</sup> 09 	1.0000F00 1.0001E01 9.95%F00 1.582F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00 1.0000F00	2.0000200

### TABLE 2

#### ROSENBROCKS FUNCTION

NSTEP	P	NFUNC	F	X+	
0	٥	3	2.4200201	1.2000500	1.0000200
1	2	13	3.6298200	<b>¯</b> ⊌.9067 <i>E</i> _01	7.69815_01
2	2	21	3.4233E00	<b>-</b> 8.4879 <i>E</i> 01	7.2767E_01
3	2	37	1.8095£00	3.3952 <i>E</i> 01	1.27608_01
Ē.	2	49	1.6006200	2.5919 <i>E</i> 01	7.9437 <i>E</i> _02
5	2	57	1.1819200	<b>7.8469</b> <i>E</i> 02	1.98658_02
6	2	65	8.8453 <i>E</i> 01	6.04628_02	7.89615_03
7	2	73	7.17385 01	1.5304 <i>E</i> _01	2.4091 <i>E</i> _02
8	2	81	5.6197 <i>E</i> 01	2.5041 <i>E</i> _01	6.17462_02
9	2	90	4.2998£ 01	3.44298 01	1.1797 <i>E</i> _01
10	2	99	3.16982 01	4.3715E 01	1.9244 <i>E</i> _01
11	2	108	2.21826 01	5.31398_01	2.87105_01
12	2	119	1.3396 <i>E</i> 01	6.34338_01	4.00802_01
13	2	127	1.0543E 01	6.7550 <i>E</i> _01	4.57448_01
14	2	140	4.4759 <i>E</i> 02	7.88468 01	6.2137E_01
15	2	150	3.0951 <i>E</i> 02	8.26775 01	6.86628_01
16	2	163	4.4350 <i>8</i> 03	9.33988_01	8.7320 <i>E</i> _01
17	2	173	2.2415E <sup>-</sup> 03	9.5434 <i>E</i> _01	9.1203 <i>2</i> _01
18	2	183	2.98015 05	9.9458E_01	9.89248 01
19	2	191	6.0225 <i>E</i> 06	9.9756 <i>E</i> _01	9.9515E_01
20	2	199	4.2134 <i>E</i> 09	9.99948 01	9.99878_01
21	1	208	8.00285 11	9.9999 <i>E</i> 01	9.9998 <i>E</i> 01
	COUV	ERGED	_		
21	1	208	8.0028E <sup>-</sup> 11	9.99995-01	9.9998 <i>E</i> 01
GZ					
5.510	⊎E_ีย	-2.77	345-8		
aa					

798.76 -401.73 -401.73 202.7

### BEALES FUNCTION

BEA	LES	FUNCTI	08		
NSTEP	P	NFUNC	F	X+	
0	0	3	1.4203201	0.0000200	0.0000E00
1	2	25	3.4995 <i>E</i> 01	2.1250E00	2.08915 01
2	2	33	8.71302 02	2.4689800	3.2620E <sup>-</sup> 01
3	2	48	5.3580 <i>E</i> 04	2.9462500	4.8493E 01
4	2	57	1.28985 04	2.9728E00	4.9262E 01
5	2	65	9.61465 07	2.9976200	4.99392 01
6	2	71	2.2789E 09	2.9999500	4.9997E 01
7	2	77	9.0499E <sup>-</sup> 13	3.0000E00	5.0000E 01
C	ONV	ERGED			
7	2	77	9.04998 13	3.0000E00	5.0000E 01
GZ					
6.826	5 <i>E</i> -	6 1.69	972E 6		
GG					
3	. 94	48 <sup>-</sup> 3	12.984		

			•	٠	٠	
- 1	2.	984	4	5		54

#### POWELLS FUNCTION 2

NSTEF	P P	NFURC	F	X+		
0	0	4	-1.5000E00	0.0000E00	1.0000E00	2.0000 <i>E</i> 00
1	3	20	-2.4976E00	2.19585 01	8.6335 <i>E</i> 01	1.6225E00
2	3	33	2.8698E00	4.0748E 01	7.73885 01	1.1885200
3	3	42	2.8974E00	4.78805 01	8.00915 01	1.1542500
4	3	57	2.9999200	1.0192200	1.0079200	9.9288 <i>E</i> 01
5	3	74	-3.0000E00	1.0051E00	1.0023200	9.98485 01
Ğ	3	85	-3.0000E00	1.0044E00	1.0017200	9.9850E 01
*IC 6						
7	3	96	3.0000E00	1.0033200	1.0019200	9.98765 01
8	3	108	3.0000E00	1.0034500	1.0015E00	9.99002 01
*IG 8	1 <sup>–</sup>	••••				
9	3	119	-3.0000E00	1.0026E00	1.0012200	9.9936E 01
10	3	133	3.0000E00	1.0000200	1.0001E00	1.0000E00
11	3	144	-3.0000E00	1.0001E00	1.0000E00	1.0000E 00
+IG 1	1	••••				
12	3	153	-3.0000E0C	1.0000E00	1.0000E00	1.0000E 00
13	3	175	3.0000E00	1.000CE00	1.0000200	9.9999 <i>E</i> 01
*SPF	13					
	CONV	ERGED				
13	3	175	-3.0000E00	1.0000500	1.0000E00	9.999 <i>9E</i> 01
2.084	8 <i>E</i> 5	1.710	5E-5 -3.17	85 <i>E</i> <sup>-</sup> 5		
GG						
4	.130	2 -	8.0504	1.8054		
- 8	.050	- -	9.0137	-0.57887		
1	. 805	÷ -	0.57887	7.7914		

### TABLE 2 (continued)

#### POWELLS FUNCTION 1

$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	1.0000E00 7.0824E 01 3.0313E 01	
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	7.0824E 01 3.0313E 01	
2 4 44 5.3349E 01 5.4315E 01 $6.4441E$ 03 3.3092E 01 3 4 64 1.9547E 01 1.4400E 01 $1.3570E$ 02 3.1944E 01 *IG 3 4 43 6.2625E 02 2.1017E 01 $1.4541E$ 02 2.3993E 01 *IG 4 4 3 6.2625E 02 1.5177E 01 $1.9142E$ 02 1.9390E 01 5 4 100 2.9855E 02 1.5177E 01 $1.9142E$ 02 1.9380E 01 4IG 5 6 4 115 1.4494E 02 1.3994E 01 $1.2674E$ 02 1.5501E 01	3.0313E 01	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	3 LO1LP 01	
• IG 3 ↓ 4 83 6.2625E <sup>-</sup> 02 2.1017E <sup>-</sup> 01 <sup>-</sup> 1.8541E <sup>-</sup> 02 2.3993E <sup>-</sup> 01 • IG ↓ 5 4 100 2.9855E <sup>-</sup> 02 1.5177E <sup>-</sup> 01 <sup>-</sup> 1.9142E <sup>-</sup> 02 1.9380E <sup>-</sup> 01 • IG 5 6 4 115 1.4494E <sup>-</sup> 02 1.3998E <sup>-</sup> 01 <sup>-</sup> 1.2674E <sup>-</sup> 02 1.5501E <sup>-</sup> 01	2142148 01	
4         4         83         6.2625E <sup>*</sup> 02         2.1017E <sup>*</sup> 01         1.4541E <sup>*</sup> 02         2.3993E <sup>*</sup> 01           FTC         5         4         100         2.9855E <sup>*</sup> 02         1.5177E <sup>*</sup> 01         1.9142E <sup>*</sup> 02         1.9380E <sup>*</sup> 01           FTC         6         4         115         1.4494E <sup>*</sup> 02         1.3998E <sup>*</sup> 01         1.2674E <sup>*</sup> 02         1.5501E <sup>*</sup> 01		
•IG 4 5 4 100 2.9#55E <sup>-</sup> 02 1.5177E <sup>-</sup> 01 <sup>-</sup> 1.9142E <sup>-</sup> 02 1.93#0E <sup>-</sup> 01 •IG 5 6 4 115 1.4494E <sup>-</sup> 02 1.399#E <sup>-</sup> 01 <sup>-</sup> 1.2674E <sup>-</sup> 02 1.5501E <sup>-</sup> 01 •IC 6	2.4747E 01	
5 4 100 2.9855E <sup>7</sup> 02 1.5177E <sup>7</sup> 01 <sup>7</sup> 1.9142E <sup>7</sup> 02 1.9380E <sup>7</sup> 01 +TC 5 6 4 115 1.4494E <sup>7</sup> 02 1.3998E <sup>7</sup> 01 <sup>7</sup> 1.2674E <sup>7</sup> 02 1.5501E <sup>7</sup> 01		
•IG 5 6 4 115 1.4494E 02 1.3998E 01 1.2674E 02 1.5501E 01	2.0664E 01	
6 4 115 1.4494E 02 1.3998E 01 1.2674E 02 1.5501E 01	_	
+TC 5	1.81265 01	
-20 0		
7 4 131 1.1635502 1.3174501 1.1921502 1.5335501	1,6854E 01	
*IG 7		
¥ 4 146 7.8424E <sup>-</sup> 03 1.2730E <sup>-</sup> 01 <sup>-</sup> 1.2215E <sup>-</sup> 02 1.4251E <sup>-</sup> 01	1.4106E 01	
+IG 8	_	
9 4 164 5.2363E <sup>-</sup> 03 1.1401E <sup>-</sup> 01 1.2318E <sup>-</sup> 02 1.2767E <sup>-</sup> 01	1.2962E <sup>-</sup> 01	
*IG 9	_	
10 4 184 4.34175 04 1.18785 01 1.14835 02 4.96085 02	5.3873E 02	
+IG 10	<b>-</b>	
11 4 203 2.91056 04 1.10956 01 1.10666 02 5.12636 02	5,2515E 02	
•IG 11		
12 4 217 2.7144E 04 1.0805E 01 1.0775E 02 5.2268E 02	5.2864E 02	
*IG 12		
13  4  233  2.6617E  04  1.0789E  01  1.0772E  02  5.1243E  02	5.2414E_02	
$14   4   248   2.5260E_04   1.0601E_01   1.0577E_02   5.1421E_02$	5.2296E 02	
<b>15 4 265 2.3636</b> <i>E</i> <b>04 1.0319</b> <i>E</i> <b>01 1.0284</b> <i>E</i> <b>02 5.1608</b> <i>E</i> <b>02</b>	5.1714 <i>E</i> -02	
*10 15		
16 4 279 2.3047E 04 1.0303E 01 1.0259E 02 5.0944E 02	5.1592E-02	
*10 16	<b>-</b> -	
17 4 294 2.25435 04 1.0257E 01 1.0235E 02 5.0364E 02	5.1364E 02	
18 4 319 2.0639E 04 1.0171E 01 1.0145E 02 4.5911E 02	4.7415E 02	
13 4 333 1.03302 04 9.87762 02 9.85702 03 4.59142 02	+.0+J32 02	
-10 L 353 1.7072FOL 0 7474FTA2 TO 6105FTA2 1 5545FTAA	L 63195 00	
eTG 20	03436 02	
21 L 372 1.71595 0L 9.61055 02 9.75995 03 L.58625 02	L.5903E 02	
*IG 21		
22 5 389 1.65532 05 9.58752 9.56362 03 5.59022 02	5793E 02	
+1G 22		
23 4 404 1.60765 04 9.53885 02 9.51885 03 4.47015 02	4.5438E.02	
24 4 420 1.22695 04 8.95295 02 8.94215 03 3.93035 02	4.0342E 02	
*IG 24		
25 4 438 1.0511E <sup>-</sup> 04 8.6293E <sup>-</sup> 02 8.6240E <sup>-</sup> 03 3.8419E <sup>-</sup> 02	3.90325 02	
<i>◆IG</i> 25		
26 4 453 1.0091E 04 8.5387E 02 8.5260E 03 3.8288E 02	3.8825502	
27 4 468 9.6162E 05 8.4156E 02 8.4060E 03 3.8470E 02	3.9061E 02	
28 4 484 6.5522E 05 7.1816E 02 7.1699E 03 3.9382E 02	3.9515E <sup>-</sup> 02	
29 4 499 5.70958 05 6.27008 02 6.26118 03 3.98008 02	3.98425 02	
30 4 513 5.60548 05 6.03448 02 5.98488 03 3.98588 02	3.97445 02	
<b>31</b> 4 526 5,3594E <sup>-</sup> 05 5,9125E <sup>-</sup> 02 <sup>-</sup> 5,9026E <sup>-</sup> 03 3,9437E <sup>-</sup> 02	3.9761E <sup>-</sup> 02	
32 4 542 5.0555E_05 5.9026E_02 5.9094E_03 3.8668E_02	3.9105E_02	
33 4 560 4.70462 05 5.88012 02 5.88142 03 3.80312 02	3.82115 02	
	3.62015 02	
34 4 579 3.8032E 05 5.6198E 02 5.6099E 03 3.5961E 02	3.1497E 02	
34 4 579 3.80328 <sup>-</sup> 05 5.61988 <sup>-</sup> 02 <sup>-</sup> 5.60998 <sup>-</sup> 03 3.59618 <sup>-</sup> 02 35 4 594 2.32218 <sup>-</sup> 05 5.03668 <sup>-</sup> 02 <sup>-</sup> 5.02898 <sup>-</sup> 03 3.08448 <sup>-</sup> 02	1.95982 02	
34 4 579 3.00322 <sup>-</sup> 05 5.6198 <sup>2</sup> 02 <sup>-</sup> 5.6098 <sup>2</sup> 03 3.5961 <sup>2</sup> 02 35 4 594 2.3221 <sup>2</sup> 05 5.0366 <sup>2</sup> 02 <sup>-</sup> 5.0298 <sup>2</sup> 03 3.0844 <sup>2</sup> 02 36 4 610 4.0940 <sup>2</sup> 06 3.6028 <sup>2</sup> 02 <sup>-</sup> 3.6038 <sup>2</sup> 03 1.9613 <sup>2</sup> 02		
34         4         579         3.80328°05         5.61988°02         5.60998°03         3.59618°02           35         4         594         2.32218°05         5.03668°02         5.02898°03         3.08448°02           36         4         610         4.0946°06         3.60288°02         3.60138°03         1.96138°02           *IG         36         -         -         -         -         -	· · · · ·	
34         4         579         3.0032E <sup>7</sup> 05         5.6198E <sup>7</sup> 02         5.6098E <sup>7</sup> 03         3.5961E <sup>7</sup> 02           35         4         5.94         2.3221E <sup>7</sup> 05         5.0366E <sup>7</sup> 02         5.0298E <sup>7</sup> 03         3.0844E <sup>7</sup> 02           4         6.10         4.0940 <sup>7</sup> 06         3.6028E <sup>7</sup> 02         3.6013E <sup>7</sup> 03         1.9613E <sup>7</sup> 02           4/6         36         -         -         -         -         -         -           37         3         6.33         4.05608 <sup>7</sup> 06         3.60407 <sup>7</sup> 02         3.6026E <sup>7</sup> 03         1.9532E <sup>7</sup> 02	1.9563 <i>E</i> 02	
34         4         579         3.8032E <sup>-</sup> 05         5.6198E <sup>-</sup> 02         5.6099E <sup>-</sup> 03         3.5961E <sup>-</sup> 02           35         4         594         2.3221E <sup>-</sup> 05         5.0366E <sup>-</sup> 02         5.0299E <sup>-</sup> 03         3.0844E <sup>-</sup> 02           36         4         610         4.0940E <sup>-</sup> 06         3.6028E <sup>-</sup> 02         3.0013E <sup>-</sup> 03         1.9613E <sup>-</sup> 02           4IG         36         4.0560E <sup>-</sup> 06         3.6040E <sup>-</sup> 02         3.6026E <sup>-</sup> 03         1.9532E <sup>-</sup> 02           4IG         37         3         633         4.0560E <sup>-</sup> 06         3.6040E <sup>-</sup> 02         3.6026E <sup>-</sup> 03         1.9532E <sup>-</sup> 02	1.9563 <i>E</i> 02	
34         4         579         3.0032£"05         5.6198£"02         5.6098£"03         3.5051£"02           35         4         5.94         2.3221£"05         5.0366£"02         5.0298£"03         3.0844£"02           36         4         610         4.0940£"06         3.6028£"02         3.6013£"03         1.9613£"02           *IG         36         3         4.0560£"06         3.6040£"02         3.6028£"03         1.9532E"02           *IG         36         3         4.0560£"06         3.6040£"02         3.6026£"03         1.9532E"02           *IG         37         3         6.30         4.0560£"06         3.6598£"02         3.5322E"02           *IG         37         4         650         3.6604£"06         3.5698£"02         3.5732£"03         1.8702E"02	1.9563 <i>E</i> <sup>-</sup> 02 1.8759 <i>E</i> -02	
34         4         579         3.00322°05         5.6198E°02         5.6098E°03         3.5961E°02           35         4         594         2.3221E°05         5.0366E°02         5.0299E°03         3.0941E°02           36         4         610         4.0940E°06         3.6028E°02         5.0299E°03         3.9941E°02           *IG         36         4         610         4.0940E°06         3.6028E°02         3.6032E°03         1.9613E°02           *IG         36         4         650         5.6040E°06         3.6040E°02         3.6026E°03         1.9532E°02           *IG         37         3         650         3.66046°06         3.5698E°02         3.5732E°03         1.9532E°02           *IG         37         4         650         3.6604E°06         3.5698E°02         3.5732E°03         1.8702E°02           38         4         655         2.5530E°06         3.3991E°02         3.3982E°03         1.4181E°02	1.9563E <sup>-</sup> 02 1.8759E <sup>-</sup> 02 1.4209E <sup>-</sup> 02	
34         4         579         3.8032E <sup>-</sup> 05         5.6194E <sup>-</sup> 02         5.6099E <sup>-</sup> 03         3.5961E <sup>-</sup> 02           35         4         594         2.3221E <sup>-</sup> 05         5.0366E <sup>-</sup> 02         5.0299E <sup>-</sup> 03         3.0444E <sup>-</sup> 02           36         4         610         4.0940E <sup>-</sup> 06         3.6024E <sup>-</sup> 02         3.6013E <sup>-</sup> 03         1.9613E <sup>-</sup> 02           4G         36         4.0560E <sup>-</sup> 06         3.6040E <sup>-</sup> 02         3.6026E <sup>-</sup> 03         1.9532E <sup>-</sup> 02           4G         37         4         650         3.6604E <sup>-</sup> 06         3.5991E <sup>-</sup> 02         3.3952E <sup>-</sup> 03           38         4         650         3.6604E <sup>-</sup> 06         3.3991E <sup>-</sup> 02         3.3935E <sup>-</sup> 03         1.4181E <sup>-</sup> 02           39         4         655         2.5530E <sup>-</sup> 06         3.3991E <sup>-</sup> 02         3.3985E <sup>-</sup> 03         1.4181E <sup>-</sup> 02           40         4         678         2.3721E <sup>-</sup> 06         3.2601E <sup>-</sup> 02         3.2598E <sup>-</sup> 03         1.4181E <sup>-</sup> 02	1.9563E <sup>-</sup> 02 1.8759E <sup>-</sup> 02 1.4209E <sup>-</sup> 02 1.1897E <sup>-</sup> 02	
34         4         579         3.00322"05         5.619£"02         5.609£"03         3.501£"02           35         4         594         2.3221£"05         5.0366£"02         5.029£"03         3.0844£"02           36         4         610         4.03402"06         3.60282"02         3.60132"03         1.9613£"02           *IG         36         4         610         4.03402"06         3.60282"02         3.60136"06         1.9532E"02           *IG         36         4         0.5608"06         3.60407"02         3.60262"03         1.9532E"02           *IG         37         3         633         4.05608"06         3.56948"02         3.732E"03         1.4702E"02           38         4         650         3.66042"06         3.56948"02         3.39945"03         1.4148"02           40         4         650         3.2504"06         3.3991E"02         3.25948"03         1.1485E"02           40         4         652         2.37218"06         3.26016"02         3.25948"03         1.1485E"02           40         4         630         1.48742"06         2.9048"02         3.25948"03         1.78398"03	1.9563E <sup>-</sup> 02 1.8759E <sup>-</sup> 02 1.4209E <sup>-</sup> 02 1.1897E <sup>-</sup> 02 8.8118E <sup>-</sup> 03	
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	1.9563E <sup>-</sup> 02 1.8759E <sup>-</sup> 02 1.4209E <sup>-</sup> 02 1.1897E <sup>-</sup> 02 8.8118E <sup>-</sup> 03 2.7756E <sup>-</sup> 03	
34         4         579         3.0032£"05         5.619£"02         5.609£"03         3.501£"02           35         4         594         2.3221£"05         5.0366E"02         5.0289£"03         3.0844E"02           36         4         610         4.09402"06         3.60285"02         5.0289£"03         3.0844E"02           416         36         -         610         -         0.9285"02         3.60126"03         1.9613£"02           416         36         -         -         0.9608"06         3.60407"02         3.6026E"03         1.9532E"02           417         3         4         650         3.66045"06         3.59915"02         3.59325"03         1.4112"02           38         4         650         3.66045"06         3.39915"02         3.25985"03         1.4112"02           39         4         665         2.55305"06         3.39915"02         3.25985"03         1.4112"02           40         4         705         1.96147"6"06         2.9048"6"02         2.9348"6"03         1.1885"02           41         4         690         1.96147"6"06         2.9048"6"02         2.9048"6"03         4.7539E"03           42         4         705         1.5	1.9563E <sup>•</sup> 02 1.8759E <sup>-</sup> 02 1.4209E <sup>-</sup> 02 1.1897E <sup>-</sup> 02 8.8118E <sup>-</sup> 03 2.7756E <sup>-</sup> 03 6.2068E <sup>-</sup> 03	
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	1.9563E <sup>•</sup> 02 1.8759E <sup>-</sup> 02 1.4209E <sup>-</sup> 02 1.1897E <sup>-</sup> 02 8.8118E <sup>-</sup> 03 2.7756E <sup>-</sup> 03 6.2068E <sup>-</sup> 03 8.2748E <sup>-</sup> 03	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1.9563E 02 1.4759E 02 1.4209E 02 1.197E 02 8.8118E 03 2.7756E 03 6.2068E 03 8.2748E 03 8.2748E 03 8.3057E 03	
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-}02\\ 1.8759E^{-}02\\ 1.4209E^{-}02\\ 1.1897E^{-}02\\ 8.8118E^{-}03\\ 2.7756E^{-}03\\ 6.2068E^{-}03\\ 8.2748E^{-}03\\ 8.3057E^{-}03\\ 8.3057E^{-}03\\ 8.368E^{-}03\\ \end{array}$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$1.9563E^{\circ}02$ $1.4759E^{\circ}02$ $1.4209E^{\circ}02$ $1.1497E^{\circ}02$ $4.4114E^{\circ}03$ $2.7756E^{\circ}03$ $4.2744E^{\circ}03$ $4.2744E^{\circ}03$ $4.3057E^{\circ}03$ $4.7649E^{\circ}03$ $9.1040E^{\circ}03$	
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$1.9563E^{-}02$ $1.4759E^{-}02$ $1.4209E^{-}02$ $1.4897E^{-}02$ $2.7756E^{-}03$ $6.2064E^{-}03$ $8.2744E^{-}03$ $8.3057E^{-}03$ $8.7649E^{-}03$ $9.1040E^{-}03$ $9.2441E^{-}03$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$1.9563E^{\circ}02$ $1.4759E^{\circ}02$ $1.49759E^{\circ}02$ $1.1497E^{\circ}02$ $4.8114E^{\circ}03$ $2.7756E^{\circ}03$ $4.2744E^{\circ}03$ $4.2744E^{\circ}03$ $4.3057E^{\circ}03$ $9.1040E^{\circ}03$ $9.241E^{\circ}03$ $9.241E^{\circ}03$ $9.241E^{\circ}03$ $9.241E^{\circ}03$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-}02\\ 1.4759E^{-}02\\ 1.4209E^{-}02\\ 1.497E^{-}02\\ 2.7756E^{-}03\\ 2.7756E^{-}03\\ 3.2745E^{-}03\\ 3.2745E^{-}03\\ 3.2744E^{-}03\\ $	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1.9563E 02 1.4759E 02 1.4209E 02 1.1497E 02 4.4114E 03 2.7756E 03 4.2744E 03 4.3057E 03 9.1040E 03 9.241E 03 9.241E 03 9.3051E 03 9.3051E 03 9.3051E 03 9.3052E 03	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-0}2\\ 1.4759E^{-0}2\\ 1.4209E^{-0}2\\ 4.4114E^{-0}3\\ 2.7756E^{-0}3\\ 2.7756E^{-0}3\\ 3.2756E^{-0}3\\ 3.2756E^{-0}3\\ 3.2748E^{-0}3\\ 3.2748E^{-0}3\\ 3.2748E^{-0}3\\ 3.2748E^{-0}3\\ 3.2748E^{-0}3\\ 3.2748E^{-0}3\\ 3.2748E^{-0}3\\ 3.288E^{-0}3\\ 3.288E^{$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-0}2\\ 1.4759E^{-0}2\\ 1.4209E^{-0}2\\ 1.1497E^{-0}2\\ 4.1148E^{-0}3\\ 2.7756E^{-0}3\\ 3.27756E^{-0}3\\ 3.2775E^{-0}3\\ 3.057E^{-0}3\\ 3.057E^{-0}3\\ 3.051E^{-0}3\\ 9.1040E^{-0}3\\ 9.3009E^{-0}3\\ 9.3009E^{-0}3\\ 9.2441E^{-0}3\\ 9.3009E^{-0}3\\ 9.2952E^{-0}3\\ 9.2441E^{-0}3\\ 9.2952E^{-0}3\\ 9.2441E^{-0}3\\ 9.2952E^{-0}3\\ 9.2441E^{-0}3\\ 9.2952E^{-0}3\\ 9.2441E^{-0}3\\ 9.2952E^{-0}3\\ 9.2443E^{-0}3\\ 9.2743E^{-0}3\\ 9.2742E^{-0}3\\ 9$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-0}2\\ 1.4759E^{-0}2\\ 1.4204E^{-0}2\\ 4.1497E^{-0}2\\ 4.1148E^{-0}3\\ 2.7756E^{-0}3\\ 5.2064E^{-0}3\\ 4.3057E^{-0}3\\ 8.3057E^{-0}3\\ 8.3057E^{-0}3\\ 9.1040E^{-0}3\\ 9.3051E^{-0}3\\ 9.3051E^{-0}3\\ 9.3051E^{-0}3\\ 9.3051E^{-0}3\\ 9.2952E^{-0}3\\ 9.2952E^{-0}3\\$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-0}2\\ 1.4759E^{-0}2\\ 1.4209E^{-0}2\\ 1.1997E^{-0}2\\ 4.4114E^{-0}3\\ 2.7756E^{-0}3\\ 3.27756E^{-0}3\\ 3.2775E^{-0}3\\ 3.0512E^{-0}3\\ 3.010+0E^{-0}3\\ 3.010+0E^{-0}3\\ 3.010+0E^{-0}3\\ 3.009E^{-0}3\\ 3.009E^{-0}3\\ 3.2441E^{-0}3\\ 3.009E^{-0}3\\ 3.2441E^{-0}3\\ 3.009E^{-0}3\\ 3.2441E^{-0}3\\ 3.2441E^{-0}3\\$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.4209E^{-02}\\ 3.1209E^{-02}\\ 4.114E^{-03}\\ 3.755E^{-03}\\ 3.755E^{-03}\\ 3.755E^{-03}\\ 3.057E^{-03}\\ 3.057$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.4209E^{-02}\\ 2.756E^{-03}\\ 2.7756E^{-03}\\ 3.7756E^{-03}\\ 3.7756E^{-03}\\ 3.7756E^{-03}\\ 3.7756E^{-03}\\ 3.748E^{-03}\\ 3.004E^{-03}\\ 3.009E^{-03}\\ 3.009E^{-03}\\ 3.2874E^{-03}\\ 3.28$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.4209E^{-02}\\ 3.4209E^{-02}\\ 3.4209E^{-02}\\ 3.4214E^{-03}\\ 3.756E^{-03}\\ 3.756E^{-03}\\ 3.057E^{-03}\\ 3.0$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.429E^{-02}\\ 2.756E^{-03}\\ 2.7756E^{-03}\\ 3.7756E^{-03}\\ 3.7756E^{-03}\\ 3.756E^{-03}\\ 3.2746E^{-03}\\ 3.051E^{-03}\\ 3.051E^{-03}\\ 3.009E^{-03}\\ 3.2952E^{-03}\\ 3.2474E^{-03}\\ 3.274E^{-03}\\ 3.44665E^{-03}\\ 3.282E^{-03}\\ 3.28$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.4209E^{-02}\\ 1.197E^{-02}\\ 4.4114E^{-03}\\ 2.7756E^{-03}\\ 3.2756E^{-03}\\ 3.2756E^{-03}\\ 3.010+0E^{-03}\\ 3.010+0E^{-03}\\ 3.010+0E^{-03}\\ 3.010+0E^{-03}\\ 3.010+0E^{-03}\\ 3.010E^{-03}\\ 3.009E^{-03}\\ 3.000E^{-03}\\ 3.000E$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.4209E^{-02}\\ 2.4209E^{-02}\\ 3.12975E^{-02}\\ 3.1148E^{-03}\\ 3.7756E^{-03}\\ 3.27756E^{-03}\\ 3.2745E^{-03}\\ 3.057E^{-03}\\ 3.057E^{-03}\\ 3.051E^{-03}\\ 3.009E^{-03}\\ 3.2745E^{-03}\\ 3.2446E^{-03}\\ 3.2446E^{-03}\\ 3.4562E^{-03}\\ 3.4665E^{-03}\\ 3.4665E^{-03}\\ 3.4665E^{-03}\\ 3.84665E^{-03}\\ 3.84665E^{-03$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.4209E^{-02}\\ 2.4209E^{-02}\\ 3.1497E^{-02}\\ 4.114E^{-03}\\ 3.7756E^{-03}\\ 3.7756E^{-03}\\ 3.7756E^{-03}\\ 3.2748E^{-03}\\ 3.051E^{-03}\\ 3.051E^{-03}\\ 3.051E^{-03}\\ 3.09E^{-03}\\ 3.09E^{-03}\\ 3.2748E^{-03}\\ 3.2748E^{-03}\\ 3.2748E^{-03}\\ 3.2748E^{-03}\\ 3.2748E^{-03}\\ 3.2748E^{-03}\\ 3.2748E^{-03}\\ 3.2748E^{-03}\\ 3.2748E^{-03}\\ 3.2874E^{-03}\\ 3.2874E$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.4209E^{-02}\\ 3.1209E^{-02}\\ 4.114E^{-03}\\ 3.755E^{-03}\\ 3.755E^{-03}\\ 3.057E^{-03}\\ 3.057$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.429E^{-02}\\ 2.4209E^{-02}\\ 4.1148E^{-03}\\ 2.7756E^{-03}\\ 2.7756E^{-03}\\ 3.62064E^{-03}\\ 3.62064E^{-03}\\ 3.62064E^{-03}\\ 3.62064E^{-03}\\ 3.62064E^{-03}\\ 3.051E^{-03}\\ 3.009E^{-03}\\ 3.092E^{-03}\\ 3.2852E^{-03}\\ 3.2474E^{-03}\\ 3.2474E^{-03}$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.429F^{-02}\\ 3.4209E^{-02}\\ 3.1497E^{-02}\\ 3.1497E^{-02}\\ 3.114E^{-03}\\ 3.057E^{-03}\\ 3.05$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1.9563E 02 1.4759E 02 1.497E 02 4.1148E 03 2.7756E 03 4.2748E 03 4.2748E 03 8.2748E 03 9.1040E 03 9.1040E 03 9.2952E 03 9.2952E 03 9.2952E 03 9.2952E 03 9.2952E 03 9.2446E 03 9.2446E 03 9.2446E 03 9.2446E 03 9.3576E 03 9.45652E 03 9.46655E 03	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1.9563E^{-02}\\ 1.4759E^{-02}\\ 1.4209E^{-02}\\ 1.1497E^{-02}\\ 4.114E^{-03}\\ 2.7756E^{-03}\\ 3.2756E^{-03}\\ 3.2756E^{-03}\\ 3.2756E^{-03}\\ 3.2756E^{-03}\\ 3.2756E^{-03}\\ 3.2748E^{-03}\\ 3.051E^{-03}\\ 3.009E^{-03}\\ 3.009E^{-03}\\ 3.092E^{-03}\\ 3.274E^{-03}\\ 9.2456E^{-03}\\ 3.274E^{-03}\\ 3.274E^{-03}\\ 3.274E^{-03}\\ 3.274E^{-03}\\ 3.274E^{-03}\\ 3.45665E^{-03}\\ 3.4665E^{-03}\\ 3.28665E^{-03}\\ 3.28665E^{$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1.9563E 02 1.4759E 02 1.497E 02 4.4114E 03 2.7756E 03 6.2064E 03 4.2744E 03 4.2744E 03 9.2454E 03 9.2454E 03 9.2454E 03 9.2454E 03 9.2454E 03 9.2454E 03 9.2454E 03 8.4562E 03 8.4665E 03 8.4665E 03	

### TABLE 3

(	CUBE						
***	rp .p	NFIL	ic )	•	X+		
0			3 7.4	04 60 2	1.2000F00	1.0000 <b>F</b> 00	
1	2	1	1.2	587E00	_1.2067E_0	1 3.505# <u>8</u> 0	3
2 • TC	2 <sup>2</sup>	2	6 1.1	8202.00	8.25176 0.	2 9.5268A U	3
-10	<b>2</b>	33	1 1.0	345800	-1.56478-02	2 5.4047 <i>E</i> 0	3
	2	4	2 5.3	925_01	2.7137E_01	1 2.0427 <b>5</b> 0	2
5	2	5	3 4.1	518E 01	3.55675 0	1 4.45936 0	2
5	2	7	2 3.1	0845 01	5.09255 01	1 1.32205 0	1
	2		1 1.8	17E 01	5.7253E 01	1.897850	ī
9	2	90	0 1.3	738E 01	6.3374E 0	1 2.60225 0	1
10	2	10:	2 8.6	5096 02	7.05836 0	1 3.50792 0	1
12	2	120	5 4.4	5148 02	7.89035 0	4.9146E 0	ī
13	2	13	3.4	41E 02	8.1536E_0	L 5.41368 <u>0</u> 0	1
14	2	14	2.4	578E 02	8.4331E 01	L 5.9923E <sup>-</sup> 0	1
15	2	15	8 1.6	1416 02 1646 02	8.73378 01 N 99778 01	1 7.28308 0	1
17	2	17	8 5.8	232 03	9.23245 01	7.8687E 0	1
18	2	1.01	3.1	316 03	9.44215 01	8.4175E 0	1
19	2	191	1.4	425 03	9.62395 01	U.9137E <sup>-</sup> 0	1
20	2	201	5.0	5088 04	9.77505 0	9.34026 0	1
21	2	221	/ 1.2 5 1.5	146E 05	9.96115 01	9.88375 0	1
23	2	23	4.8	50E 07	9.99315 01	9.9792E 0	1
24	2	243	2 1.9	537E 09	9.99965 01	9.9987 <u>8</u> 0	1
25	1	254	5.2	9868-15	1.0000E 00	1.00004 0	0
25	201	251	5.2	N65-15	1.00005-00	1.0000E 0	0
	72 <sup>-</sup>		• ••••				-
2.94	99E	¥ _8'	. #3428 9				
ć	7 <b>6</b>						
		•					
	600	.0.	200.3	25			
			RUNCET	• <i>"</i>			
, na start and a start	(NULL	1810.	0.4866	2 -1.30	03 -0.89519		
					-		
NSTER	, p	NFURC	; r 5 710	2502		-1 1216500	
ĭ	3	13	9.97	0E 00	-4.9704E-01	-1.4301E00	7.7209E 01
2	3	27	8.030	60E00	<b>~.94608~01</b>	1.3932200	7.85948 01
+IG 2	2	•••					
3	3	38	3.63	9200	_4.9302E_01	1.3809400	- W.1216K 01
5	3	58	1.468	05.00	~4.9145E~01	1.3433600	8.5029E 01
6	з	68	1.107	2E_02	_4.8701 <i>E</i> _01	[1.3039 <b>E</b> 00	_8.9160E_01
?	3	77	4.921	2E 03	-4.86885-01	1.3027E00	_w.9272E_01
	3	105	1.00/	35 05	4.86685 01	1.3008200	- 8.94715 01 - # 9516F 01
10	3	115	1.283	45 08	-4.8662E-01	-1.3004E00	-#.9519E-01
11	1	130	1.571	8E 09	4.8662E 01	1.3003500	8.95192 01
●SPP	11						
11	CONVE	130	1.674	HE-09		-1.3003800	
î cz	•	135				2.3003200	4.33136 01
-0.00	07675	ig 0.	00072181	0.0004	.7837		
GG							
30	505.7	,	21567	153	17		
-	891.9	9	15387	1309	1		

## TABLE 3 (continued)

STEP	P NFU	NC F	1+		
0	0	2.5000E03	1.0000E00	0.0000500	0.0000200
1	3 2	8 2.2316E01	9.5036E 01	6.767 HE 01	4.1635200
2	3 3	9 1.6428501	B.352#E 01	6.23825 01	6.0168800
IG 2					
з	3 4	9 1.523#E01	7.8195E 01	6.37908 01	3.9016800
4	36	0 1.3986201	6.9259E 01	7.09845 01	3.73#1200
5	3 7	7.4964200	1.5051E 01	9.75978 01	2.7332800
6	3 9	6.5887E00	5.000HE 02	1.0018500	2.5518800
7	3 11	2.7227E00	5.2671E-01	8.61358 01	1.6409200
8	3 12	5 2.5228E00	5.54648 01	8.54878 01	1.5731200
IGB			••••		
9	3 14:	3 2.1290E00	6.39258 01	7.97598 01	1.1370800
10	3 154	1.7566E00	6.722#8 01	7.40125 01	1.3253800
11	3 16	1.3967E90	7.36035 01	6.53355 01	1.1567200
12	3 19	4.9120E 01	9.1941E 01	4.2433E 01	6.8933E
13	3 201	2.1926E <sup>-</sup> 01	9.60115 01	2.8562E 01	4.6523E 0
14	3 21	7 1.2787E <sup>-</sup> 01	9.81895 01	2.1753E 01	3.36235 0
15	3 235	7.3891 <i>E</i> 02	9.89755 01	1.68975 01	2.68745 0
16	3 240	1.6776E 02	9.96735 01	8.16955 02	1.290450
17	3 251	9.94795 03	9.98165 01	6.2932E 02	9.90535 0
18	3 261	3.6942E 03	9.9912E <sup>-</sup> 01	3.82645 02	6.0732E 0
19	3 280	1.3546 <i>E</i> 03	9.99848 01	2.2736E 02	3.6579E 0
20	3 291	3.82235 04	9.93935 01	1.23205 02	1.95385 0
21	3 304	6.4540E 05	9.99995 01	5.05925 03	8.0306E 0
22	3 313	9.3914E 06	1.0000 <i>E</i> 00	1.91528 03	3.0616E 0
23	3 325	3.0591E <sup>-</sup> 06	1.0000E 00	1.10445 03	1.74145 0
24	3 335	9.53425 07	1.00005 00	6.1509E 04	9.75988 0
25	3 346	2.94518 07	1.0000E00	3.39885 04	5.4246E 0
26	3 356	4.5557 <i>E</i> 08	1.0000E00	1.34735 04	2.128850
27	3 369	6.48425 09	1.0000E 00	5.0713E 05	8.0496E 0
28	3 376	1.1147E 09	1.0000E 00	2.09595 05	3.338620
29	2 394	3.57195 10	1.0000E 00	1.19345 05	1.88125 0
30	1 412	8.2403E 12	1.00005 00	1.78812 06	2.83805 0
C	ONVERGEL	)			
30	0 424	7.9056E 12	1.0000 <i>E</i> 00	1.72485 06	2.7451E 0
GZ	_				
. 4 2 4 9	57 1.1	199E 5 6.7893	E 6		

200.04	118.89	71.555
118.89	520.95	322.86
71.555	322.86	202.25

WOODS FUNCTION

NSTEP	P	NFUNC	,	X+			
0	0	5	1.9192504	3.0000E00	1.0000E00	<b>3.0000E00</b>	1.000 <i>0E</i> 00
ī	ĩ.	26	3.6916501	1.38535 01	3.6686E 01	2.2318E 01	1.92265 01
2	÷.	52	1.0818201	2.0079E 01	1.25925 01	1.1184E00	1.1970200
ATC 2	-	•-					
	L	71	#.339#E00	3.6696501	1.86998 01	-1.1208E00	1.2626800
•16 3	-	• •			•••••		
-10 0	L	35	7.7625800	L.3L2HE 01	2.25025 01	-1.1071E00	1.2945800
-TG -	-	•••					
-10 -	L		7.0012800	5.09878-01	2.57825 01	1,1190200	1.3027500
	7	116	6.27#2/00	5 23568 01	3.00965 01	1.1571800	1.3530800
	7	110	6 9495 <b>7</b> 00	5.13342 01	3 17615 01	1.177#800	1.3892200
	۰.	130	3.3893500	3.41132 01	3.1.010 01		
-16 /				6 7010F-01	3 22835-01	-1 1762E00	1 6075200
		144	3.8308200	5./2126 01	3.11936 01		1
•16 •			· · · · · Boo	· · · · · · · · · · · · · · · · · · ·	3 3503F 01		1
	•	123	5./411600	2.04835 01	3.33028 01	1.18302.00	1.4130500
•16 9				c	3 61 10F 01		1 5922800
10	. *	179	2.4203800	P.03337 01	3.34108 01	1.2333200	1.5422800
+IG 10				· · · · · · · · · · · · · · · · · · ·	3 6066E-04		4 61 33 844
11		193	5.4006200	6.00362 01	3.69665 01	-1.23/0800	1.5422600
12		206	5.3824200	6.06182 01	3.73432 01	1.2358200	1.538/800
13	4	223	5.3669500	6.13598 01	3.79108 01	1.2386200	1.5365800
14	4	237	5.3471500	6.1642E 01	3.85862 01	1.2331200	1.5300500
+IG 14			· · · · · · · · · · · · · · · · · · ·		<b></b>		
15	4	252	5.3216 <i>B</i> 00	6.275#E_01	3.9249E_01	1.2312200	1.5211200
16	4	269	5.279#£00	6.4026 <u>8</u> 01	4.12798_01	1.2324800	1,5231800
17	4	286	5.2274800	6.71528_01	4.4786E_01	1.2351200	1.5274200
18	4	298	5.2080200	6.74828_01	4.56755_01	1.2308200	1.5243200
19	4	317	4.9765200	7.23728-01	5.28788 01	1.1497200	1.3258500
20	4	341	3.9000 <i>E</i> 00	1.0574 <i>E</i> 00	1.1229500	_7.7958E_01	6.5765 <u>8</u> 01
21	4	360	3.6285500	1.1004 <i>E</i> 00	1.2100500	_7.3603 <i>E</i> _01	5.5792E_01
22	4	377	3.253#200	1.1631800	1.3806800	_6.5712 <i>E</i> _01	4.4784 <b>8</b> _01
23	4	389	3.126 <i>8E</i> 00	1.1777E00	1.395 <i>8E</i> 00	_6.9536E_01	4.9994 <i>8</i> _01
24	4	401	3.0794500	1.1905E00	1.4191200	_6.9126 <i>E</i> _01	4.90518_01
25	4	417	2.9547200	1.2413E00	1.5458800	_6.6256 <i>E</i> _01	4.4919E_01
26	4	435	2.8254800	1.2839800	1.6382200	_6.0128E_01	3.71448_01
27	4	452	2.5623200	1.3573E00	1.8338200	~4.5276E~01	2.22045 01
28	4	468	2.5010E00	1.3441800	1.8076200	~~.57~5E~01	2.08655 01
29	÷.	486	1.9926200	1.3835500	1.915#200	2.1881E 01	5.60122 02
30	Ĺ.	510	1.7757800	1.3980200	1.9570500	1.0740E 01	2.24945 02
31	Ĩ.	526	1.5361200	1.41#2500	2.0041500	2.31785 02	3.9660E 03
32	ĩ.	539	1.3258200	1.4139200	1,9999200	1.3001E 01	2.10815 02
11	ĩ	555	1.1672800	1.4104500	1.9909500	2.36576 01	5.2709E 02
34	7	576	1.1332500	1.6111200	1.9813200	2.46265 01	5.61858 02
~ ~	-						

## TABLE 3 (continued)

• I G	34															
35		4	596	1	.02	77 <b>EO</b> O		1	. 383	7 <b>E</b> O	0	1.	916	2500	2.81518_01	6.5163E <sup>-</sup> 02
36		4	613		. 24	55 <i>2</i> - 0	1	1	. 337	850	0	1.	791	6 <i>E</i> 00	3.7191 <i>E</i> 01	1.49465 01
37		4	632	7	. 59:	16 <b>2</b> 0	1	1	. 296	450	0	1.	677	7500	4.4681E <sup>-</sup> 01	2.01385 01
3.4		É.	649	6	. 79	55 0	1	1	271	120	٥	1.	634	6800	4.97815 01	2.53758 01
			671	ĩ	6.7			-	1 97		<u> </u>	4.	LL7	*800	6. LH1HE 01	L 17648 01
		7			••••			- 11			~				6 34705 01	
40				3					• • • • •	320				3200	6.36/08 01	4.151/2 01
41		4	700	3	. 35	665 0	1	1	.247	61.0	0	1.	224	7800	6.40138 01	4.08948 01
42		4	715	3	.27	50E 0	1	1	.246	920	0	1.	556	5500	6.4635 <b>8</b> 01	4.13518 01
+IG	42															
6.3		ι.	729	3	. 22'	7750	1	1.	. 246	# <i>E</i> 0	0	1.	557	00.00	6.4623E 01	4.1788E 01
+ TG	1.3															
			71.6						21.6		•	4.		#200	6 6221 8 01	L 27748-01
		2						- 11		350	~		61.7	0500	6 00245 01	1 71605 01
			/63		• / 4	002 0				100				5200	6.90212 01	4.74005 01
46		4	779	2	. 69	887.0	1	1.	. 240	920	0	1.	236	62.00	6.89218 01	4.7337E 01
47		4	794	2	.67	-1 <i>E</i> _0	1	1.	. 239	1 <i>E</i> 0	0	1.	535	1500	6.8879 <i>E</i> _01	4.7470E_01
48		4	809	2	. 55	∎5 <i>E</i> 0	1	1.	. 232	6 <i>E</i> 0	0	1.	519	9200	6.9361 <i>E</i> 01	4.8172E 01
49		4	#23	2	.25	242 0	1	1.	. 222	380	0	1.	494	8200	7.1959E 01	5.1652E 01
50		É.		1		11510	1	1.	1 # 2	950	٥	1.	396	3500	7.99625 01	6.37645 01
		7								250	~			6 200	H 13035 01	C 61767 01
31										250	~		340	1500	0.13025 01	6.61/08 01
52			884		. 23	- 22 0	1	1.	.1/4	36 0		1.	303		0.07592 01	6.50502 01
53		4	896	1.	. 22	58E. 0	1	1.	.170	1 E O	0	1.	370	00.00	8.05652 01	6.47428 01
♦IG	53					_									-	-
54		4	913	1	.21:	0 <i>-</i> 300	1	1.	.170	2E0	٥	1.	367	2800	<b>8.0594</b> <i>E</i> 01	6.49215 01
+IG	54															
55	•	μ.	929	1.	. 201	97 0	1	1.	168	920	0	1.	366	9500	8.05935 01	6.4795E 01
. 16		•		•			-				-				•••••••••••••••••••••••••••••••••••••••	
											•		166		N A61 67 A1	A 10337-01
30		•	344		• 1 3 :	302 0					•	* •	303	3800	0.00435 01	0.49310 01
+1G	56															<b></b>
57		4	959	1	.180	<b>545 0</b>	1	1.	167	3 <i>E</i> 0	0	1.	363	35.00	8.07362 01	6.50412 01
+IG	57					_									_	_
58		4	977	1.	.17:	10 <i>E</i> 0	1	1.	.167	3 E 0	0	1.	361	0500	8.0857E 01	6.5203E 01
+IG	5 #															
60		L	001	1	. 1 1 0	5 F 0	1	1	161	1	•	1.	151	3500	8.12795-01	6 6029F 01
					••••		:	- 11		250	~			5 500		6 61 1 F 01
			100/		••••									3500		6.04445 01
61		•	1023			22_0	2	1	134	120	0	1.	283	/200	8.33302 01	0.94302 01
62		4	1036	7.	. 602	21E 0	2	1.	.114	3 <i>E</i> 0	0	1.	249	1 <i>E</i> .00	8.5059E 01	7.2376E 01
•IG	62					_										_
63		4	1052	5.	. 6 9 2	≥9 <i>E</i> 0	2	1.	.116	620	٥	1.	247	1800	8.6497E <sup>01</sup>	7.48325 01
61		μ.	1065	5.	. 319	3E 0	2	1.	. 1 1 1	250	0	1.	235	2E00	N.6917E 01	7.56858 01
		É.	1078		66.6	25-0	2	-	005	OF O	- -	- E	10.	5500	# #152F 01	7 77685 01
		7	1006				-	- 11		350			202	1.500	9 076CF 01	
		7	1030				:				×	•••	200	3500	N 00775 01	w 00235 01
		•	1111	3.		0 310	4			120		••	200	3200	0.39/25 VI	•
+1G	67														· · · · · · · •	
68		4	1131	3,	. 507	75E_0	2	1.	. 0 9 5	OEO	0	1.	193	51.00	9.00288_01	8.09905_01
69		4	1160	3.	.13	75 <i>E</i> 0	2	1.		2 <i>E</i> 0	0	1.	184	4 <i>E</i> 00	9.0142E <sup>-</sup> 01	¥.1318E <sup>-</sup> 01
70		4	1175	2.	. 320	50 <i>E</i> 0	2	1.	065	6 <i>E</i> 0	٥	1.	135	2E00	9.18985 01	8.44968 01
71		É.	1193	1.	.379	6E 0	2	1	0.00	CEO	0	1.	124	6200	9.38518 01	H. HO45E 01
72		Ĺ.	1207		. 51	95-0	ā	1	016	1 5 0	ň	- E .	093	8800	9.60615 01	9.23178 01
		7	1 2 2 4				š			720			000	2800	9 55775 01	A 12528 01
/3			1441				3			/80				2500	9.33372 01	9.12328 01
74		•	1237	<b>D</b> .	. 875	198_0	3	1.	043	250	0	1.	088	4200	9.5/5/2 01	9.16668 01
75		4	1257	5.	. 996	<u>-</u>	3	1.	038	850	0	1.	079	35.00	w.5788E_01	9.1758E 01
76		4	1271	3.	.10	94E_0	3	1.	.029	2 E O	D	1.	059	2 <i>E</i> 00	9.70558_01	9.4190E_01
77		4	1288	3.	. 581	76E 0	4	1.	.007	4 E O	D	1.	014	9 <i>E</i> 0 0	9.94915 01	9.89045 01
78		4	1301	3.	. 804	.9 <i>E</i> 0	5	1.	003	2 2 0	0	1.	900	3E00	9.9671E <sup>-</sup> 01	9.9343E 01
70		ĩ.	1315		.03	0E-0	ŝ	1	002	920	0	1	005	8800	9.97305 01	9.94595 01
		7					2							7500		0 071 45 01
			1330				2			35.0		**			8.30305 01	
81		•	1344	1.	. 8 4 3	112_0	6	- y .		<u></u> `	1	a.	338	. 01	1.0003200	1.0005800
82		4	1356	2.	. 207	"4E_0	8	9.	. 999	8E_'	01.	9.	999	68 01	1.0000200	1.0001200
83		4	1378	1.	.162	22E_0	8	9.	. 999	7 E 🗍	01	9.	999	3 <i>E</i> _01	1.000 <i>0E</i> 00	1.0001500
84		1	1396	6.	.134	•BE_0	9	9.	. 999	6 <i>E</i> -	01	9.	999	2E-01	1.0000200	1.0001700
85		1	1415	2.	. 529	5E 0	9	9.	999	8E -	01	9.	999	7 <i>E</i> 01	1.0000500	1.0000 <i>E</i> 00
86		ī	16.12		07	6E -	9		990	9.5	01	9.	990	85 01	1.0000500	1.0000700
	~	- •	2434	•			-				•			**		
							•	~				•				1
		•	1454		• 936	100 1	v	у,	993	AC	11	ж.		05 01	T.0000V00	1.0000200
86		0			_				-			-				
¥6	z,	0					A4 A6 4	17	- L	. 91'	79E	5				
86 -2.5	72 3491	0 #2 <sup></sup> 5	2.0	5154E	5	0.00	01001	•••				-				
86 -2.3	72 94 91	0 #Z <sup>-</sup> 5	2.0	5154E	- 5	0.00	01083	•••		•••		-				
86 -2.3	72 74 91 76	0 #Z <sup>-</sup> 5	2.0	5154E	- 5	0.00	0106	•••	•	•••		-				
86 -2.3	72 94 91 76	0 #Z <sup>-</sup> 5	2.0	5154E	- 5	0.00	01083	•••	•			•				
86 -2.3	72 3491 76 78	0 #2 <sup>-</sup> 5	2.1	5154E	-5	0.00	51	. 1 1	2		1.1	. 93	8			
86 -2.3	72 1491 76	0 #Z <sup>-</sup> 5	5 2.0	-380.	-5 .23	0.00	_51	. 1 1	2		18	.93	8			
86 -2.: (	72 3491 76 789	0 #2 <sup>-5</sup> 9.6 0.23	2.0	-380. 205.	.23	0.00	-51 -23	. 1 1	2	-	18	. 93	5 7			
86 -2.: (	72 3491 76 789 380 51	0 82 <sup>-</sup> 5 0.6 0.23 1.11	2.0	-380 205 23	5 .23 .12 .93	0.00		. 11	2	-	18	.93 .51 .26	5 7			

(Figures taken from [14])						
Method	QNWD	H-J	Ros	SPLX	Pow	Stew
Function						
Rosenbrock	208(-11)	250(-8)	200(-6)	200(-8)	151(-10)*	163(-12)**
Beale	77(-13)	$100(-\infty)$	130(-7)	100(-8)		
Powell 1	978(-7)		• •		433(-13)*	407(-10)**
Cube	254(-15)		200(-∞)	140(-7)		
Box	191(-11)	100(-∞)		290(-5)		
RTF(3)***	130-284				96-120	
	Av. = 189				Av. = 108	
RTF(5)***	312-406				166 — 167	
	Av. = 370				Av. = 166	

1	<b>FABLI</b>	e 4	
Comparison	with	Other	Methods

These figures come from [10].

\*\* These figures come from [4].

\*\*\* These are Random Trigonometric Functions of dimension 3 and 5. The accuracy criterion used is that the maximum error in any x-component is  $<10^{-7}$ . The smallest and largest numbers of evaluations taken are listed, as well as the averages.

down to a certain value is listed. The number in parentheses is the exponent, to base 10, of the least calculated function value. The value " $-\infty$ " indicates that f was reduced to zero.

7. Acknowledgments. I am especially indebted to Dr. P. G. Comba, whose suggestions and criticisms sowed the seeds for many of the ideas in this work. I am also grateful to Drs. D. Goldfarb and Y. Bard for very helpful discussions, and to Jean-Claude Cohen for his help in setting up the program.

Thanks are also due to M. J. D. Powell whose criticisms resulted in a simplification and clarification of the presentation.

The bulk of this work was done at the IBM New York Scientific Center.

### Appendix A.

Limiting Cases of  $\nu \to 0$  and  $\nu \to \infty$ .

Case 1:  $\nu \to 0$ . If  $\nu$  is set to zero in Eq. (4.17a), the formula for  $\theta_1$  is not defined, since  $\tau_1^2 = \sigma_1^2$ . Therefore, we must consider  $\theta_1$  (and  $\eta_1$ ) separately. The formula for  $\theta_1$  is:

$$\theta_1 = (\epsilon_1 - 2\rho_1)/\sigma_1^2$$

and for  $\eta_1$ , we have:

(A2) 
$$\eta_1 = 4\nu^2 \rho_1 / \sigma_1^4 - 2\theta_1.$$

When  $i \neq 1$ , we have:

(A3) 
$$\theta_i = \frac{2\nu^2(\epsilon_i - 2\rho_i)}{\sigma_i^2(\tau_i^2 - \sigma_i^2)} + O(\nu^4),$$

JOHN GREENSTADT

(A4) 
$$\eta_i = \frac{4\nu^2}{\sigma_1^4} \rho_i - 2\theta_i.$$

When  $\nu \to 0$ , every term in formula (4.5a) goes to zero, except the first term (for i = 1). The result for  $\gamma$  is (also replacing  $\Lambda$  by L):

(A5) 
$$\gamma \to L\theta_1 \sigma_1.$$

For  $\Gamma$ , we must be more careful. When we replace M by  $L/\nu$ , we have a denominator which converges to 0, whereas  $\theta_1$  and  $\eta_1$  do not. However, if we evaluate the terms in the brace in formula (4.5b) for i = 1, we obtain:

(A6) 
$$\{ \}_{i=1} = \eta_1 \sigma_1 \sigma_1^T + 2\theta_1 \sigma_1 \sigma_1^T$$

since  $\tau_1 = \sigma_1$ . Replacing  $\eta_1$  by expression (A2), we then have:

(A7) 
$$\{ \}_{i=1} = \frac{4\nu^2}{\sigma_1^4} \rho_1 \sigma_1 \sigma_1^T - 2\theta_1 \sigma_1 \sigma_1^T + 2\theta_1 \sigma_1 \sigma_1^T ,$$

so that all we have left is the first term. There is no difficulty with the rest of the terms in Eq. (4.5b).

For convenience, we define:

(A8) 
$$\tilde{\theta}_i \equiv \frac{2(\epsilon_i - 2\rho_i)}{\sigma_i^2(\tau_i^2 - \sigma_i^2)}; \quad i \neq 1,$$

(A9) 
$$\tilde{\eta}_1 \equiv \frac{4\rho_1}{\sigma_1^4} ,$$

(A10) 
$$\tilde{\eta}_i \equiv \frac{4\rho_i}{\sigma_i^4} - 2\tilde{\theta}_i; \quad i \neq 1,$$

so that

$$\theta_i/\nu^2 = \tilde{\theta}_i + O(\nu^2)$$
 and  $\eta_i/\nu^2 = \tilde{\eta}_i + O(\nu^2)$ .

Then  $\Gamma$  becomes (replacing *M* by  $L/\nu$ ):

(A11)  

$$\Gamma = \frac{1}{2\nu^2} L \left\{ \frac{4\nu^2 \rho_1 \sigma_1 \sigma_1^T}{\sigma_1^4} + \sum_{i \neq 1} \left[ \eta_i \sigma_i \sigma_i^T + \theta_i (\sigma_i \tau_i^T + \tau_i \sigma_i^T) \right] \right\} L$$

$$= \frac{1}{2} L \left\{ \tilde{\eta}_1 \sigma_1 \sigma_1^T + \sum_{i \neq 1} \left[ \tilde{\eta}_i \sigma_i \sigma_i^T + \tilde{\theta}_i (\sigma_i \tau_i^T + \tau_i \sigma_i^T) \right] \right\} L + O(\nu^2)$$

and when  $\nu \rightarrow 0$ , the last term vanishes.

Clearly, this limiting procedure has the effect of correcting  $g_0$  from the results of the first minor step only, and of removing part of the first minor step discrepancy from the correction to G.

Case 2:  $\nu \to \infty$ . In this case, there is no need to separate out the first minor step. The limit for  $\theta_i$  is:

(A12) 
$$\theta_i \to (\epsilon_i - 2\rho_i)/\sigma_i^2$$
,

but  $\eta_i$  still contains a multiple of  $\nu^2$ . The formula for  $\gamma$  remains the same as (4.5a), but that for  $\Gamma$  becomes:

(A13) 
$$\Gamma = \frac{1}{2} L \left\{ \sum_{i} \frac{4\rho_i}{\sigma_i^4} \sigma_i \sigma_i^T \right\} L + O\left(\frac{1}{\nu^2}\right)$$

and the last term vanishes for  $\nu \to \infty$ . In this case,  $g_0$  is corrected in terms of all the minor steps, but the G-correction does not contain the  $\theta$ 's.

In the program used to run the test problems, L was set equal to the unit matrix I, as mentioned in the text.

#### Appendix B.

Comparison with Fiacco-McCormick Method. The method described by Fiacco and McCormick in their book [16] is based largely on a relation identical with Eq. (3.15a). Let a step  $\sigma$  be made up of a linear combination of at most two coordinate directions, viz.:

(B1) 
$$\sigma_{(ij)} = \alpha_i e_i + \alpha_j e_j.$$

That is, let the direction  $S_{ij}$  be specified in terms of coordinate directions  $e_i$  and  $e_j$ , and do a line search for the minimum of f along that direction, starting at a point  $x_0$ . Then the minimum is found at  $x_1 (\equiv x_0 + \sigma_{(ij)})$  and the difference between starting and minimum values of f is denoted by  $\Delta f_{(ij)}$ . We then have, rewriting (3.15a):

(B2) 
$$\Delta f_{(ij)} = -\frac{1}{2}\sigma_{(ij)}^T G^* \sigma_{(ij)}$$

and, replacing  $\sigma_{(ij)}$  according to (B1), we obtain:

(B3) 
$$\Delta f_{(ij)} = -\frac{1}{2} \{ \alpha_i^2 e_i^T G^* e_i + 2\alpha_i \alpha_j e_i^T G^* e_j + \alpha_j^2 e_i^T G^* e_j \}$$

(remembering that  $G^*$  is symmetric). But, because the coordinate-direction vector  $e_i$  has the structure:  $e_i = (0, 0, \dots, 0, 1, 0, \dots, 0)$ —where the 1 is in the *i*th position—each of the products singles out a component of  $G^*$ . Thus, e.g.,

$$(B4) e_i^T G^* e_i = G_{ii}^*$$

so that (B3) becomes:

(B5) 
$$\Delta f_{(ij)} = -\frac{1}{2} \{ \alpha_i^2 G_{ij}^* + 2\alpha_i \alpha_j G_{ij}^* + \alpha_j^2 G_{jj}^* \}$$

Now, we choose the first set of directions for  $\sigma$  so that they lie along the coordinates. Then, we have:

$$\Delta f_{(ii)} = -\frac{1}{2}\alpha_i^2 G_{ii}^*,$$

from which we can solve for the diagonal elements  $G_{ii}^*$ .

Next, we arrange that  $\alpha_i = \alpha_i$  (and denote them both by  $\alpha_{ii}$ ), i.e., we search in a direction (always starting at  $x_0$ , as before) which bisects the right angle between  $e_i$  and  $e_i$ . We then have:

(B7) 
$$\Delta f_{(ij)} = -\frac{1}{2}\alpha_{ij}^2 (G_{ii}^* + G_{ij}^* + 2G_{ij}^*),$$

from which we can solve for  $G_{ij}^*$ , since everything else is known. Clearly, since  $G^*$  is symmetric, we need only have done  $\frac{1}{2} N(N + 1)$  line searches.

Once we have estimated  $G^*$  in this way, we make use of Eq. (3.15b), using the results of the searches along the coordinate directions. ( $\tau$  is, of course, the same as  $\sigma$  for a single line search.) We then have:

 $\alpha_i g_{0i}^* + \alpha_i^2 G_{ii}^* = 0,$ (B8)

from which we solve for  $\{g_{0,i}^{*}\}$ . We may then translate  $g^{*}$  to any other point, using (3.12).

The main differences between this method and the QN method outlined in this paper are:

F-M	QN

1.	$\frac{1}{2} N(N + 1)$ line searches	N line searches
2.	Complete estimate of $g_0$ and $G$	Incomplete estimate of $g_0$ and $G$
	(exact for a quadratic function)	
3.	Completely new estimate at next	Improvement of previous estimates at
	major step	next major step

**IBM Scientific Center** 2670 Hanover Street Palo Alto, California 94304

1. W. C. DAVIDON, Variable Metric Method for Minimization, Argonne National Lab., ANL-5990 Rev., 1959.

ANL-3990 Rev., 1939.
2. R. FLETCHER & M. J. D. POWELL, "A rapidly convergent descent method for minimization," Comput. J., v. 6, 1963/64, pp. 163-168. MR 27 #2096.
3. C. G. BROYDEN, "Quasi-Newton methods and their application to function minimization," Math. Comp., v. 21, 1967, pp. 368-381. MR 36 #7317.
4. G. W. STEWART III, "A modification of Davidon's minimization method to accept difference approximations of derivatives," J. Assoc. Comput. Mach., v. 14, 1967, pp. 72-83. MR 35 #7566.

5. D. GOLDFARB, "Sufficient conditions for the convergence of a variable-metric algorithm," in Optimization, R. Fletcher (Editor), Academic Press, New York, 1969.
6. G. P. MCCORMICK & J. D. PEARSON, "Variable metric methods and unconstrained optimization," in Optimization, R. Fletcher (Editor), Academic Press, New York, 1969.
7. J. GREENSTADT, "Variations on variable-metric methods," Math. Comp., v. 24, 1970, pp. 1-22. MR 41 #2895.

8. H. R. SOSENBROCK, "An automatic method for finding the greatest or least value of a function," Comput. J., v. 3, 1960/61, pp. 175-184. MR 24 #B2081.
9. E. M. L. BEALE, On an Iterative Method for Finding a Local Minimum of a Function of More than One Variable, Stat. Tech. Res. Grp., Technical Report #25, Princeton University Devices the Princeton of More than One Variable, Stat. Tech. Res. Grp., Technical Report #25, Princeton

University, Princeton, N. J., 1958.
10. M. J. D. POWELL, "An iterative method for finding stationary values of a function of several variables," *Comput. J.*, v. 5, 1962, pp. 147–151.
11. M. J. D. POWELL, "An efficient method for finding the minimum of a function of several variables without calculating derivatives," *Comput. J.*, v. 7, 1964, pp. 155–162. MR 32 #4828.

12. A. LEON, "A comparison among eight known optimizing procedures," in *Recent* Advances in Optimization Techniques, A. Lavi and T. P. Vogl (Editors), Wiley, New York, 1960.

13. D. GOLDFARB, "A family of variable-metric methods derived by variational means," *Math. Comp.*, v. 24, 1970, pp. 23-26. MR 41 #2896.

14. J. KOWALIK & M. R. OSBORNE, Methods for Unconstrained Optimization Problems, American Elsevier, New York, 1968.

15. A. R. COLVILLE, A Comparative Study of Non-Linear Programming Codes, IBM

Tech. Rep. #320-2949, 1968. 16. A. V. FIACCO N G. P. MCCORMICK, Nonlinear Programming: Sequential Uncon-strained Minimization Techniques, Wiley, New York, 1968, p. 175 et seq. MR 39 #5152.